

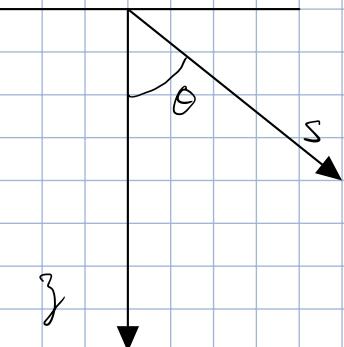
4-5pm 588 vector calc

take an

temp drop over couple cm
over little path, temp. practically is constant

WKB approximation

Radiative diffusion



at large T $I_\nu \rightarrow S_\nu'$
thermal radiation $S_0 = B_\nu(T)$

show that $F \propto \frac{dT}{dz}$

$$\mu = \cos \theta$$

$$\mu = \frac{z}{s}$$

$$ds = \frac{dz}{\cos \theta} = \frac{dz}{\mu}$$

$$\mu \frac{d}{dz} (I_\nu(z, \mu)) = -(\alpha_\nu + \sigma_\nu) (I_\nu - S_\nu')$$

$$I_\nu(z, \mu) = S_\nu - \frac{\mu}{\alpha_\nu + \sigma_\nu} \frac{dI_\nu}{dz}$$

use zeroth order as first order correction

$$\text{b/c } I_\nu^{(0)}(z, \mu) = S_\nu^{(0)}(z, \mu) . \quad J_\nu^{(0)} = S_\nu^{(0)} \Rightarrow I_\nu^{(0)} = S_\nu^{(0)} = B_\nu$$

$$I_\nu = I_\nu^{(0)} + I_\nu^{(1)} + \dots$$

$$I_{\nu}^{(1)} = I_{\nu}^{(0)} - \frac{1}{\alpha_{\nu} + \beta_{\nu}} \cdot \frac{\partial}{\partial z} (B_{\nu})$$

treat $I_{\nu} = I_{\nu}^{(0)}$

now compute flux

$$F_{\nu}(z) = \int I_{\nu}^{(1)}(z, \mu) \cos \theta d\Omega$$

$$= 2\pi \int_{-1}^{+1} I_{\nu}^{(1)}(z, \mu) \cdot \mu d\mu$$

bc B_{ν} is isotropic, it has no angular dependencies.
Integrating over all angles will make it zero. can neglect

$$F_{\nu}(z) = -\frac{2\pi}{\alpha_{\nu} + \beta_{\nu}} \frac{\partial B_{\nu}}{\partial z} \cdot \int_{-1}^{+1} \mu^2 d\mu$$

$$= -\frac{4\pi}{3(\alpha_{\nu} + \beta_{\nu})} \frac{\partial B_{\nu}(T)}{\partial T} \cdot \frac{dT}{dz}$$

$$B_{\nu} = B(\nu, T)$$

to get bolometric flux, integrate over all frequencies

$$F(z) = \int_0^{\infty} F_{\nu}(z) d\nu$$

$$= -\frac{4\pi}{3} \frac{dT}{dz} \cdot \int_0^{\infty} \frac{1}{(\alpha_{\nu} + \beta_{\nu})} \cdot \frac{\partial B_{\nu}}{\partial T} d\nu$$

in photon
flux of heat
not radiated

Stefan
Bo耳emann

$$F(z) = \frac{-160T^3}{3\alpha_R} \frac{dT}{dz}$$

$$\int_0^{\infty} \frac{\partial B_{\nu}}{\partial T} d\nu = \frac{4\pi T^3}{\pi}$$

$$\alpha_R^{-1} = \frac{\int_0^{\infty} (\alpha_{\nu} + \beta_{\nu})^{-1} \cdot \frac{\partial B_{\nu}}{\partial T} d\nu}{\int_0^{\infty} \frac{\partial B_{\nu}}{\partial T} d\nu}$$

↑
Rosseland absorption

all assuming temp. kinda stays same
≠ very optically thick

Maxwell Equations

Lorentz Force: $\vec{F} = q(\vec{E} + \frac{\vec{v}}{c} \times \vec{B})$

gives meaning to electric field proportion constant introduced w/o speed of light mentioned.

admit charge for electric field like admit mass for gravitational field

Rate of doing work

$$\vec{v} \cdot \vec{F} = q \vec{v} \cdot \vec{E} \quad (\vec{B} \text{ does no work})$$

for non-relativistic particles: $\vec{F} = m \frac{d\vec{v}}{dt}$

$$\frac{d}{dt} \left(\frac{1}{2} m \vec{v}^2 \right) = q \vec{E} \cdot \vec{v}$$

electric field is only field to change KE

generalize to distribution of charges & currents

$$\vec{F} = \rho \vec{E} + \frac{1}{c} \vec{j} \times \vec{B}$$

force per unit volume
charge density
 $\rho = \int f(x) \delta(x - x_0)$
current density
 $j = \int f(x) dx$

$$\int f(x') \delta(x - x') dx' = f(x)$$

corresponding eq for work done energy/volume

$$\frac{d}{dt} (U_{\text{mean}}) = \vec{j} \cdot \vec{E}$$

Evolution of \vec{E} & \vec{B} governed by Maxwell's eqn

$$\vec{\nabla} \cdot \vec{D} = 4\pi\rho$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \times \vec{H} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{D}}{\partial t}$$

$$\vec{D} = \epsilon \vec{E}$$

$$\vec{H} = \mu \vec{B}$$

↳ displacement current predicts waves

in vacuum: $\epsilon = \mu = 1$

depends on units

$$\frac{\partial}{\partial t} (\rho) + \vec{\nabla} \cdot \vec{F} = 0$$

$$\left[\frac{\text{stuff}}{\text{volume}} \right] + \left[\frac{\text{stuff}}{\text{area} \cdot \text{time}} \right] = 0 \rightarrow \text{conservation of stuff}$$

$$\vec{\nabla} \times \vec{F} = 0 \rightarrow \vec{F} = \vec{\nabla} \phi \quad \text{rotation}$$

$$\vec{\nabla} \cdot \vec{F} = 0 \rightarrow \vec{F} = \vec{\nabla} \times \vec{A} \quad \text{solid/ideal}$$

$$\vec{F} = \vec{F}_I + \vec{F}_S$$