

$$P^\mu = m_0 U^\mu$$

$$U^\nu U_\mu = -c^2$$

$$P^\mu P_\mu \rightarrow c^2 = m_0^2 c^4 + c^2 (\vec{p})^2$$

$$P^\mu = \begin{pmatrix} E/c \\ \vec{p} \end{pmatrix} = \begin{pmatrix} \gamma c \\ \vec{\gamma} \vec{p} \end{pmatrix}$$

$$\text{define 4 acceleration : } a^\mu = \frac{dU^\mu}{dt}$$

$$P^\mu P_\mu = 0 \quad F^\mu = m_0 a^\mu = \frac{dP^\mu}{dt}$$

For EM should be related to Lorentz Force ($F_{\mu\nu} U^\nu$)

$$F_{LF} = \frac{e}{c} F_\nu^\mu U^\nu$$

$$F_\nu^\mu = F^{\mu\nu} / (m_0 c) = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} F^{\mu\nu}$$

$$F_{\mu\nu} = A_{\nu\mu} - A_{\mu\nu}$$

$$F^{\mu\nu} = \begin{bmatrix} 0 & 1 & 1 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$$

$$\frac{e}{c} F_\nu^\mu U^\nu = (E \vec{\alpha} \cdot \vec{u}) \left(\frac{\partial c}{\partial u_\nu} \right) = \frac{e\vec{r}}{c} \left(\frac{\vec{E} \cdot \vec{u}}{e\vec{r}(\vec{E} \cdot \vec{u})} \right)$$

$$\frac{1}{m_0} \frac{d}{dt} (\vec{p}) = \frac{1}{m_0} \frac{d}{dt} \left(\frac{e\vec{r}}{c} \right)$$

$$\frac{d}{dt} \left(\frac{\vec{r}}{c} \right) = e(\vec{E} + \vec{u} \times \vec{B}) \rightarrow \frac{d\vec{r}}{dt} = (\vec{E} \cdot \vec{u}) e$$

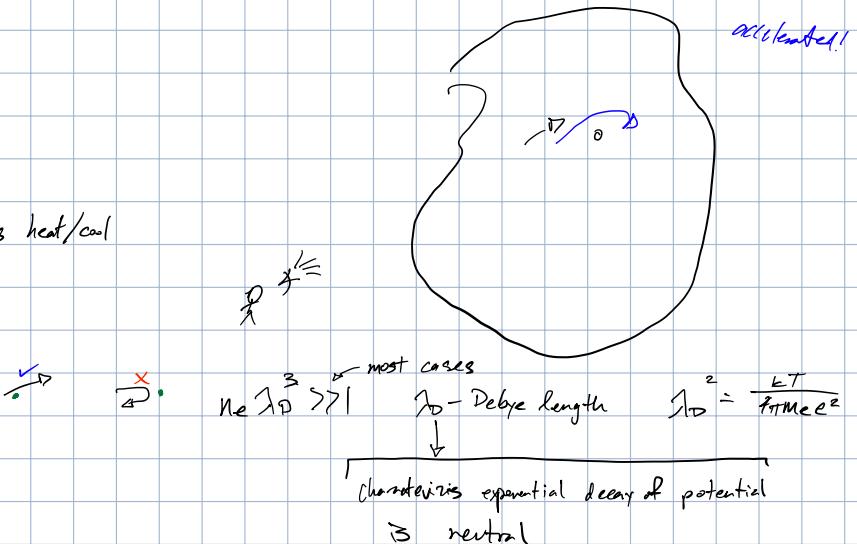
Bremsstrahlung

coolumb interactions

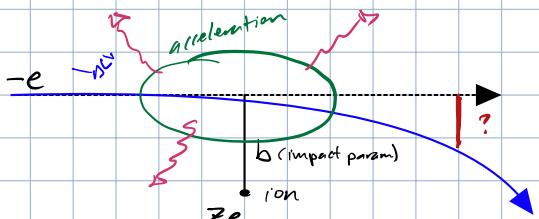
what does \vec{r} see? exactly how clouds of ionized gas heat/cool

assumptions

- ① Non adiabatic
- ② No large angle collisions
- ③ Ions are stationary

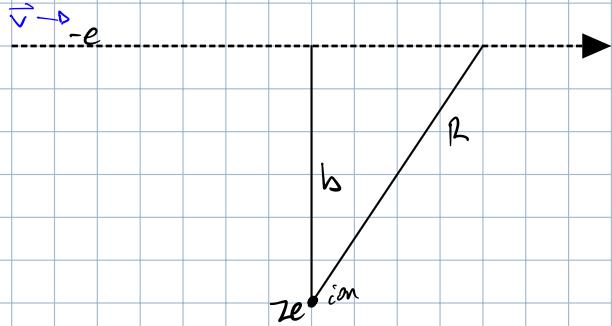


3 steps



in box $(\lambda_D)^3$ quasineutral, want halo particles to fluidize

but we have lot of nucleons. will interact with



use dipole approximation

$$\vec{d} = \sum e_i \vec{v}_i \quad \text{at origin e ion's pos}$$

$$\vec{d} = -e\vec{R} \quad \ddot{\vec{d}} = -e\ddot{\vec{v}}$$

take Fourier Transform

$$-\omega^2 \hat{d}(w) = -\frac{e}{2\pi} \int_{-\infty}^{\infty} \dot{v} e^{iwt} dt$$

particles interact strongly: $T = b\nu$ (collision temp)

for $2\omega \gg 1$, $\exp(i) \approx 1$, oscillates fast, low effect

for $2\omega \ll 1$, $\exp(i) \approx 1$

To lowest order:

$$\hat{d}(w) = \begin{cases} \frac{e}{2\pi w} \Delta V & \text{for } 2\omega \ll 1 \\ 0 & \text{for } 2\omega \gg 1 \end{cases}$$

ΔV is change in velocity due to collision

$$\frac{dw}{d\omega} = \begin{cases} \frac{2e^2}{3\pi c^2} |\Delta V|^2 & \omega \tau \ll 1 \\ 0 & \omega \tau \gg 1 \end{cases}$$

Assume ΔV is in transverse direction only

integrate transverse component of acceleration

$$\Delta V = \frac{Ze^2}{m} \int_{-\infty}^{\infty} \frac{b \cdot dt}{(b^2 + v^2 t^2)^{3/2}} \approx \int \frac{1}{b^2}$$

$$\Delta V = \frac{2Ze^2}{mbV}$$

Single collision & small angle scattering gives

$$\frac{dW(b)}{dw} = \begin{cases} \frac{8Z^2 e^6}{3\pi c^2 m^2 v^2 b^2} & b \ll v_w \\ 0 & b \gg v_w \end{cases} \quad \star$$

Suppose we have plasma

$$n_i - \text{ion density} \quad (\text{fixed speed } V)$$

$$n_e - \text{electron density}$$

electron flux incident on ion $n_e V$

area element around single ion: $2\pi b db$

emission / time / volume / frequency

$$\frac{dW}{dt dV dw} = n_e n_i 2\pi \int_{b_{\min}}^{b_{\max}} dW(b) b db = (\text{constant}) \cdot \ln\left(\frac{b_{\max}}{b_{\min}}\right) \quad \star$$

what are b_{\min} & b_{\max} ?? ? id?

b_{\max} : number of order $\frac{V}{w}$. take $w = b_{\max}$

b_{\min} : exclude particles getting too close (keep small angle approximation)

① small approximation not valid? $b_{\min}^{(1)} = \frac{4Ze^2}{\pi m v} \quad \Delta V \sim v$

② QM starts taking place

$$\Delta x = b \quad \Delta p = mv \quad \rightarrow \Delta x \Delta p \geq \hbar/2$$

$$b_{\min}^{(2)} = \frac{\hbar}{mv}$$

$$\frac{1}{2}mv^2 \ll Z^2 \left(\frac{me^4}{2t^2} \right)$$

$$\ln\left(\frac{b_{\max}}{b_{\min}}\right) \Rightarrow g(v, \omega) \quad \text{Giant Factor}$$

everything else stays