

No more codes in

Vacuum  $E \& B$ 's

$$\nabla \cdot \vec{B} = 0 \quad (1)$$
$$\nabla \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} = 0 \quad (3)$$

$$\nabla \cdot \vec{E} = 4\pi\rho \quad (2)$$
$$\nabla \times \vec{B} - \frac{1}{c} \frac{\partial \vec{E}}{\partial t} = \frac{4\pi}{c} \vec{j} \quad (4)$$

in presence of  $\rho \& \vec{j}$   
cgs units! not MKS

QM systems w/ charged particles interacting w/ classical  $E \& B$  fields

$$\nabla \cdot \vec{B} = 0 \quad \text{can be solved w/ } \vec{B} = \nabla \times \vec{A} \quad (\text{bc } \nabla \cdot (\nabla \times \vec{A}) = 0)$$

$$\rightarrow \nabla \times \left( \vec{E} + \frac{1}{c} \frac{\partial \vec{A}}{\partial t} \right) = 0$$

$$\vec{E} = -\frac{1}{c} \frac{\partial \vec{A}}{\partial t} + \text{something w/ zero curl}$$

$$\vec{E} = -\frac{1}{c} \frac{\partial \vec{A}}{\partial t} - \nabla \phi$$

$$\vec{B} = \nabla \times \vec{A} \quad (1)$$
$$\vec{E} = -\frac{1}{c} \frac{\partial \vec{A}}{\partial t} - \nabla \phi \quad (2)$$

$\vec{A}$  - vector potentials  
 $\phi$  - scalar potentials

can plug into (3) & (4) to get wave eqs  
for  $\vec{A}, \phi$

$$\vec{A}_\mu = (\phi, \vec{A})$$

$\hat{=}$  gauge field

$$j^\mu = (\rho, \vec{j})$$

$\vec{A} \& \phi$  are not unique

$$\nabla \cdot (\nabla \cdot f) = 0$$

$$\vec{A}' = \vec{A} + \nabla \lambda(\vec{x}, t)$$

$$\phi' = \phi - \frac{1}{c} \frac{\partial}{\partial t} (\lambda(\vec{x}, t))$$

} gauge transformation

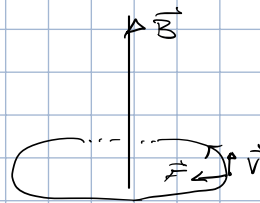
$\vec{E}, \vec{B}$  are same whether use  $(\phi, \vec{A})$  or  $(\phi', \vec{A}')$

"choose a gauge" - impose condition on  $\vec{A} \& \phi$  to remove ambiguity

$$\vec{A} = (A_x(x, y, z, t), A_y(x, y, z, t), A_z(x, y, z, t))$$

Classically, particle w/ charge  $q$  obeys equations of motion (eom)

$$m \frac{d^2 \vec{x}}{dt^2} = q \vec{E} + \frac{q}{c} \left( \frac{d\vec{x}}{dt} \times \vec{B} \right)$$



$q = -e$  for electron  $\forall e > 0$

the com follows from Hamiltonian

$$H = \frac{1}{2m} \left( \vec{p} - \frac{q}{c} \vec{A} \right)^2 + q\phi$$

$q$  - charge of charged particle

use Hamilton's eq's

$$\begin{aligned} \vec{A}(\vec{x}, t) &= \vec{A} \\ \phi(\vec{x}, t) &= \phi \end{aligned}$$

$$\vec{p} \cdot \vec{p} - \frac{q}{c} (\vec{p} \cdot \vec{A} + \vec{A} \cdot \vec{p}) + \left( \frac{q}{c} \right)^2 \vec{A} \cdot \vec{A}$$

in classical mechanics, cool

in QM: in general

$$[\vec{p}, \vec{A}(\vec{x}, t)] \neq 0$$

$$[p_i, A_j] f(\vec{x})$$

$$\begin{aligned} &= -i\hbar \frac{\partial}{\partial x_i} (A_j(\vec{x}, t) f(\vec{x})) - A_j(\vec{x}, t) (-i\hbar \frac{\partial}{\partial x_i} f(\vec{x})) \\ &= -i\hbar \frac{\partial A_j}{\partial x_i} f(\vec{x}) - i\hbar A_j \frac{\partial f(\vec{x})}{\partial x_i} + i\hbar A_j \frac{\partial f(\vec{x})}{\partial x_i} \\ &= -i\hbar \frac{\partial A_j}{\partial x_i} f(\vec{x}) \end{aligned}$$

cancel!

① Check  $H$  leads to com

② What happens when we use  $H$  in QM?

$$[p_i, A_j] = -i\hbar \frac{\partial A_j}{\partial x_i}$$

if  $\nabla A = 0 \rightarrow$  woo, this is 0!

Consider SE

$$i\hbar \frac{\partial \psi(\vec{x}, t)}{\partial t} = H \psi(\vec{x}, t) = \left( \frac{1}{2m} \left( \vec{p} - \frac{q}{c} \vec{A}(\vec{x}, t) \right)^2 + q\phi(\vec{x}, t) \right) \psi(\vec{x}, t)$$

$\vec{p} \rightarrow -i\hbar \nabla$  in pos<sup>n</sup> space       $\vec{A}$  in pos<sup>n</sup> space

This eq<sup>n</sup> is not invariant under gauge transformations of  $\vec{A}$ ,  $\phi$

same  $\vec{E} \neq \vec{B}$ , different SE

$$\text{in QM: } \psi \rightarrow \psi' = e^{i q \Lambda(x,t) / \hbar c} \psi$$

$(\vec{A}', \phi', \psi')$  has same physics as  $(\vec{A}, \phi, \psi)$

$$\begin{aligned} \text{Define } D_0 &= \frac{\partial}{\partial t} + \frac{i q \phi}{\hbar} \\ \vec{D} &= \vec{\nabla} - \frac{i q}{\hbar c} \vec{A} \end{aligned} \quad \left. \vphantom{\begin{aligned} D_0 \\ \vec{D} \end{aligned}} \right\} \text{covariant derivatives}$$

$$\text{SE: } i \hbar D_0 \psi = -\frac{\hbar^2}{2m} \vec{D} \cdot \vec{D} \psi$$

$$\begin{aligned} D_0' \psi' &= \left( \frac{\partial}{\partial t} + \frac{i q}{\hbar} \phi' \right) \left( e^{i q \Lambda / \hbar c} \psi \right) \\ &= \frac{\partial}{\partial t} \left( e^{i q \Lambda / \hbar c} \psi \right) + \frac{i q}{\hbar} e^{i q \Lambda / \hbar c} \psi \\ &= \frac{i q}{\hbar c} \frac{\partial \Lambda}{\partial t} e^{i q \Lambda / \hbar c} \psi + e^{i q \Lambda / \hbar c} \frac{\partial \psi}{\partial t} + \frac{i q}{\hbar} e^{i q \Lambda / \hbar c} \psi \\ &= \left( \frac{i q}{\hbar c} \frac{\partial \Lambda}{\partial t} + \frac{\partial}{\partial t} + \frac{i q}{\hbar} \right) e^{i q \Lambda / \hbar c} \psi \end{aligned}$$