

Fun w/ magnetic monopoles (just fun)

① Is QM consistent w/ long-range magnetic monopole fields?

$$\vec{B} = \frac{g}{r^2} \hat{e}_r \quad \nabla \cdot \vec{B} \neq 0$$

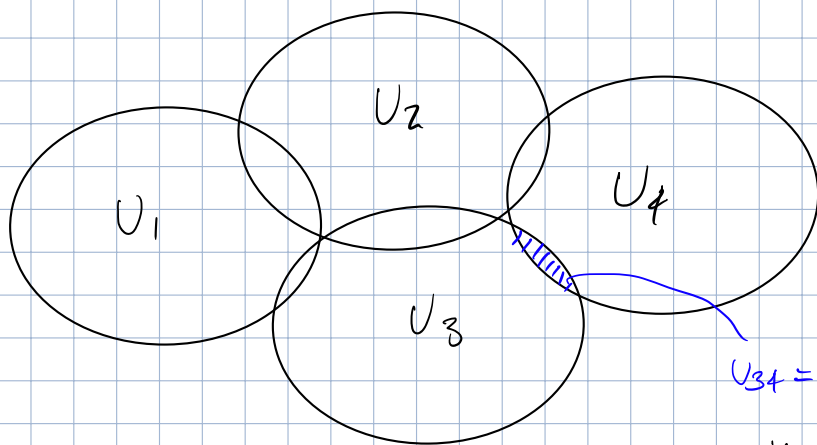
② Do actual theories predict existence of magnetic monopoles?

↳ need string theory, not today
only talk about QM

Standard Model?	No
Grand Unified Theory?	Yes
Kaluza-Klein Theory?	Yes
String Theory?	Yes

Fiber bundle

a flexible framework for gauge transformations & \vec{B} fields ($\vec{E} = 0, \Phi = 0$)
can be included



$$U_{34} = U_3 \cap U_4$$

in each U_i , will have vector potential \vec{A}_i ,
a wave $A \sim |\psi_i\rangle$, & $\vec{B}_i = \nabla \times \vec{A}_i$

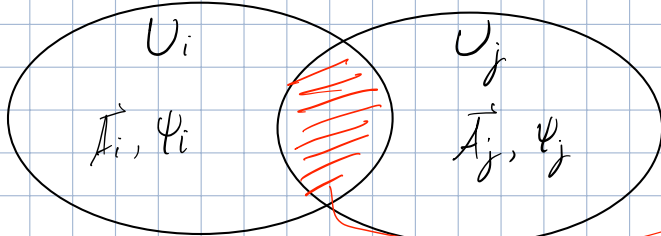
we want a well defined \vec{B} everywhere

if $U_{ij} = U_i \cap U_j$ then U_{ij} better have $\vec{B}_i = \nabla \times \vec{A}_i = \vec{B}_j = \nabla \times \vec{A}_j$

$\Rightarrow \vec{A}_i \neq \vec{A}_j$ are related by gauge transformation $\rightarrow \vec{A}_i = \vec{A}_j + \nabla \lambda_{ij}$

in QM, gauge transformations also work on wave ψ

$$\Rightarrow |\psi_i\rangle = e^{i\frac{q}{\hbar c} \lambda_{ij}} |\psi_j\rangle$$



relate i to j using gauge transformation
w/ function λ_{ij}

can rewrite $\vec{\nabla} \lambda_{ij} = \frac{\hbar c}{ig} e^{-\frac{ig\lambda_{ij}}{\hbar c}} \vec{\nabla} e^{\frac{ig\lambda_{ij}}{\hbar c}}$

$$g(\lambda_{ij}) = e^{\frac{ig\lambda_{ij}}{\hbar c}}$$

$$g^{-1}(\lambda_{ij}) = e^{-\frac{ig\lambda_{ij}}{\hbar c}}$$

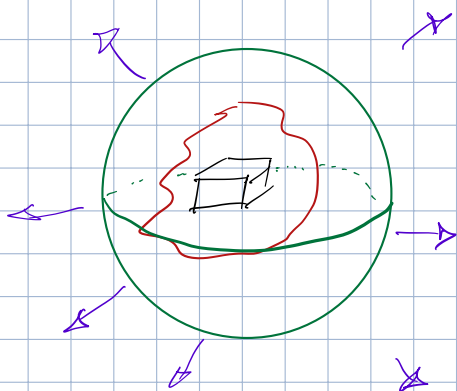
$$\psi_i = g(\lambda_{ij}) \psi_j$$

$$A_i = g^{-1}(\lambda_{ij}) \left(\vec{A}_j - \frac{\hbar c}{g} \vec{\nabla} \right) g(\lambda_{ij})$$

$$g(\lambda_1) g(\lambda_2) = g(\lambda_1 + \lambda_2)$$

$$g(\lambda) g^{-1}(\lambda) = \mathbb{1}$$

gauge group of transformation

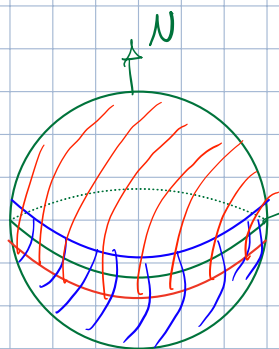


ide whats inside

can only describe \vec{B} outside

specifically, some sphere outside

$$\vec{B} = \frac{g}{r^2} \hat{e}_r$$



North & lil south U_N
 equator
 South & lil north U_S

define $\vec{A}_N \neq \psi_N$
 $\neq \vec{A}_S \neq \psi_S$

(r, θ, ϕ)

on North $0 \leq \theta \leq \frac{\pi}{2} + \epsilon$

$$\vec{A}_N = \frac{g}{r} \frac{(1 - \cos\theta)}{\sin\theta} \hat{e}_\phi$$

$$\vec{B}_N = \vec{\nabla} \times \vec{A}_N = \frac{\hat{e}_r}{r \sin\theta} \left(\frac{\partial}{\partial \theta} (\sin\theta \hat{e}_\phi) - \frac{\partial A_\phi}{\partial \phi} \right) + \cancel{\hat{e}_\theta} \left[\right] + \cancel{\hat{e}_\phi} \left[\right]$$

$$= \frac{g}{r^2} \hat{e}_r$$

\vec{A}_N is well defined on N as $\theta \rightarrow 0$
 but \vec{A}_N is singular @ $\theta = \pi$

$$\frac{1 - \cos \theta}{\sin \theta} \sim \theta \sim 0$$

on Sath

$$\vec{A}_S = -\frac{g}{r} \frac{1 + \cos \theta}{\sin \theta} \hat{e}_\theta \quad \text{well defined @ } \theta = \pi$$

compute $\vec{B} = \nabla \times \vec{A}_S = \frac{g}{r^2} \hat{e}_r$

SO $\vec{B} = \frac{g}{r^2} \hat{e}_r$ in $N \neq S$

overlap region around equator

$$\vec{A}_N = g^{-1}(\lambda_{NS}) \left(\vec{A}_S - \frac{ikc}{g} \nabla \right) g(\lambda_{NS})$$

@ $\theta = \pi/2$, $\vec{A}_N = \frac{g}{r} \hat{e}_\theta$
 $\vec{A}_S = -\frac{g}{r} \hat{e}_\theta$

$$\vec{A}_N = \vec{A}_S + \nabla \lambda_{NS} \quad \forall \lambda_{NS} = 2g \varphi$$

$$g(\lambda_{NS}) = e^{i\theta(2g\varphi)/\hbar c}$$

Is g_{NS} a well defined f^k along equator?

$g(\lambda_{NS})$ is f^k of φ

$$g(\lambda_{NS}, \varphi) = g(\lambda_{NS}, \varphi + 2\pi)$$

b/c $g(\lambda_{NS})$ is periodic f^k

$$\Rightarrow \frac{2\pi g}{\hbar c} 2\pi = 2\pi n$$

$$n \in \mathbb{Z}$$

$$\frac{n}{2} = \frac{qg}{\hbar c}$$

Dirac quantization condition

① g will be $\frac{\hbar c}{q} \cdot \frac{n}{2} = |g|$ for some n

② if there's 1 magnetic monopole anywhere, by consistency w/ QM

$$\Rightarrow q = \frac{n}{2} \frac{\hbar c}{g}$$

where g is strength of magnetic monopole

\rightarrow predicts quantization of electric charge

lower limit of g ?

n -winding# of wrap around S^1

$$q \cdot g = \frac{n}{2} \cdot \hbar c$$

"principle $U(1)$ fibre bundle"

n = value of first Chern class of this bundle