

States  $\sigma|\psi\rangle$

operators:  $A \rightarrow \sigma A \sigma^{-1}$

$$K|\psi\rangle = |\psi\rangle^* \\ K\sigma K^{-1} = \sigma^*$$

2 loose ends including spin

Time reversal  $T: \vec{L} = \vec{r} \times \vec{p} \rightarrow -\vec{L} \quad \dagger \quad \vec{S} \rightarrow \vec{S}$

$$T\vec{S}T^{-1} = -\vec{S} \quad \text{for spin } 1/2: \vec{S} = \frac{\hbar}{2} \vec{\sigma}$$

$$T = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} K$$

unitary operator  $U$  on  $\mathbb{C}^2$       complex conjugation

$$\begin{aligned} T\sigma_x T^{-1} &= U K \sigma_x K^{-1} U \\ &= \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} K \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} K^{-1} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \\ &= \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}^* \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \\ &= \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix} = -\sigma_x \end{aligned}$$

$$T\sigma_y T^{-1} = -\sigma_y$$

$$T\sigma_z T^{-1} = -\sigma_z$$

$$K(\ )K^{-1} = (\ )^*$$

$$T^2 = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} K^2 = \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix} \mathbb{1} = -\mathbb{1}$$

if  $H$  is a time reversal invariant Hamiltonian, then  $[H, T] = 0$

$$H|\psi\rangle = E|\psi\rangle \rightarrow H T|\psi\rangle = T H|\psi\rangle = T E|\psi\rangle = E \cdot T|\psi\rangle$$

$|\psi\rangle$  &  $T|\psi\rangle$  have same energy

2 possibilities:

①  $|\psi\rangle, T|\psi\rangle$  are same state  $\rightarrow T|\psi\rangle = e^{i\phi} |\psi\rangle \quad (T^2 = 1)$

②  $|\psi\rangle, T|\psi\rangle$  are distinct states w/ same energy  $\rightarrow T^2 = -1$

Okay take  $\odot$  & assume  $|X\rangle \neq$  zero vector

then  $T|X\rangle = e^{i\phi}|X\rangle$

$$\rightarrow T^2|X\rangle = T e^{i\phi}|X\rangle = \underbrace{T e^{i\phi} T^{-1} T}_{\substack{\rightarrow \text{complex conjugate} \\ \rightarrow T^{-1}T=1}}|X\rangle = e^{-i\phi} T|X\rangle = e^{-i\phi} e^{i\phi}|X\rangle = |X\rangle$$

so if  $T|X\rangle = e^{i\phi}|X\rangle$ , then  $T^2|X\rangle = |X\rangle$ , or  $T^2 = 1$

if  $T^2 = -1 \rightarrow T|X\rangle = -|X\rangle \rightarrow$  so  $|X\rangle$  is zero vector

any system  $-N$  electrons, interacting  $\vec{E} \neq 0$  &  $\vec{B} = 0$

if  $T^2 = (-1)$  (net spin  $1/2$ )

$\rightarrow$  2 fold degeneracy in energy spectrum (Kramer's Thm)

act on spin up  $T \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ -1 \end{pmatrix} = -1 \rightarrow$

What about effects of spin for electron in a  $\vec{B}$  field?

electrons have a magnetic moment:  $\vec{\mu} = g \left( \frac{-e}{2m_e c} \right) \vec{S}$

$$H_{\text{spin}} = -\vec{\mu} \cdot \vec{B} = \frac{eB}{m_e c} \left( \frac{1}{2} \right) S_z$$

spin up & spin down have different energies

$$\Delta E_{\text{st}} = \frac{eB\hbar}{m_e c} \left( \frac{1}{2} \right)$$

$$\Delta E_{\text{LL}} = E_{n+1} - E_n = \hbar \omega_B = \frac{eB\hbar}{m_e c} \text{ 'oh yes?'}$$

in vacuum,  $g \approx 2 \rightarrow \underline{\Delta E_{\text{spin}} = \Delta E_{\text{LL}}}$

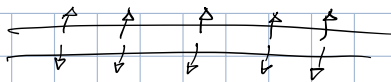
physics

MOSFET (GaAs)

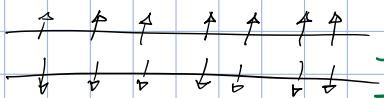
interactions change  $m_e \rightarrow m_e^{\text{eff}}$  &  $g \rightarrow g_{\text{eff}}$

$$\Delta E_{\text{spin}} \sim \frac{1}{70} \Delta E_{\text{LL}} \quad (\text{Zeeman effect})$$

\*not relevant for materials



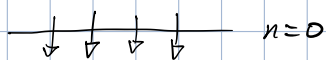
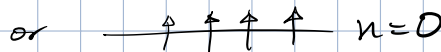
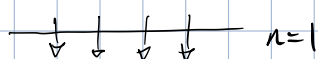
10x



$n=0$  Landau level  
lowest energy

math

$$g \equiv 2$$



2 states w/same energy  $\rightarrow$  degeneracy why?

degeneracy due to structure

$$H = \{Q, Q\} \quad [H, Q] = 0$$

$n=1$  supersymmetry

## Magnetic Monopoles

Old Maxwell's eqs  
in vacuum

$$\left. \begin{aligned} \vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} &= 0 \\ \vec{\nabla} \cdot \vec{E} &= 0 \\ \vec{\nabla} \times \vec{B} - \frac{1}{c} \frac{\partial \vec{E}}{\partial t} &= 0 \end{aligned} \right\}$$

Symmetry

$$D: \begin{aligned} \vec{E} &\rightarrow \vec{B} \\ \vec{B} &\rightarrow -\vec{E} \end{aligned}$$

duality

$$D^2: \begin{aligned} \vec{E} &\rightarrow -\vec{E} \\ \vec{B} &\rightarrow -\vec{B} \end{aligned}$$

$$\vec{\nabla} \cdot \vec{E} = 4\pi \rho_e$$

$$\vec{\nabla} \cdot \vec{B} = 4\pi \rho_m \quad (0 \text{ usually})$$

if so,  $\rho_e \rightarrow \rho_m$  &  $\rho_m \rightarrow -\rho_e$

magnetic monopoles

Dirac, Wu & Yang  $\sim 1954$

$\uparrow$

QM  $\Rightarrow$  electric charge is quantized

$$N_g = \frac{Ne}{\hbar c}$$

$\hookrightarrow N \in \mathbb{Z}$  integer