

"But in QM, things are more complicated"

Rotations in 3-D & implementation in QM

- ① Group elements depend on 3 parameters
- ② Rotation group is nonabelian

Translations: $\phi(x) \rightarrow \phi(x-e) = T(e)\phi(x) = e^{-i e P/\hbar} \phi(x)$

$\psi(\vec{x}, t) \in \mathcal{H}$
 $|\psi(x)\rangle \in \mathcal{H}$

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} \rightarrow R \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

↙ 3x3

$x^2 + y^2 + z^2$ - invariant $\rightarrow R^T R = 1$

$\phi(\vec{r}') \rightarrow \phi(\vec{r})$

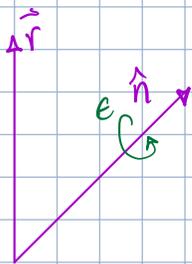
specify rotation by axis \hat{n} (leave this direction fixed) & an angle e

$\phi(\vec{r}') = e^{-ie\hat{n}\cdot\vec{L}} \phi(\vec{r})$

where $\vec{L} = \vec{r} \times \vec{p}$
 $\hat{L} = \hat{r} \times \hat{p}$

3 generators L_x, L_y, L_z

$L_i = \sum_j \sum_k \epsilon_{ijk} r_j p_k$



$\vec{r} \rightarrow \vec{r} - e \hat{n} \times \vec{r} + \mathcal{O}(e^2)$

$\vec{r} = r_n \hat{n} + \vec{r}_\perp \quad \vec{r}_\perp \cdot \hat{n} = 0$

$\hat{n} = \hat{e}_z$

$\vec{r} \rightarrow \vec{r} - e \hat{e}_z \times \vec{r} = \vec{r} - e \hat{e}_z \times (x\hat{e}_x + y\hat{e}_y + z\hat{e}_z)$

$x \rightarrow x + ey \quad z \rightarrow z$
 $y \rightarrow y - ex$

$$\begin{pmatrix} x \\ y \end{pmatrix} \rightarrow \begin{pmatrix} x' \\ y' \end{pmatrix} = \begin{pmatrix} \cos & \sin \\ -\sin & \cos \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} \stackrel{\text{small } \epsilon}{\sim} \begin{pmatrix} 1 & \epsilon \\ -\epsilon & 1 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x + \epsilon y \\ y - \epsilon x \end{pmatrix}$$

Want to show for $\hat{n} = \hat{e}_z$

$$\phi(x', y', z') = e^{-i\epsilon L_z / \hbar} \cdot \phi(x, y, z) \quad \text{to order } \epsilon$$

↑
generator of rotations about \hat{e}_z

$$\phi(x + \epsilon y, y - \epsilon x, z) = \phi(x, y, z) + \epsilon y \frac{\partial \phi}{\partial x} - \epsilon x \frac{\partial \phi}{\partial y} + \mathcal{O}(\epsilon^2) \quad \text{Taylor series}$$

Right hand side

$$\begin{aligned} e^{-i\epsilon L_z / \hbar} \phi(x, y, z) &= e^{-i\epsilon(xp_y - yp_x) / \hbar} \phi(x, y, z) \\ &= \left(1 - \frac{i\epsilon}{\hbar} (xp_y - yp_x) + \mathcal{O}(\epsilon^2) \right) \phi(x, y, z) \\ &= \left(1 + y \frac{\partial}{\partial x} - \epsilon x \frac{\partial}{\partial y} + \mathcal{O}(\epsilon^2) \right) \phi(x, y, z) \end{aligned}$$

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X

$$\begin{aligned} [L_x, L_y] &= [yp_z - zp_y, zp_x - xp_z] \\ &= yp_x [p_z, z] + p_y x [z, p_z] \\ &= -i\hbar (yp_x - p_y x) = i\hbar (xp_y - yp_x) = i\hbar L_z \end{aligned}$$

$$L_i = (L_1, L_2, L_3)$$

$$[L_i, L_j] = i\hbar \epsilon_{ijk} L_k$$

$$\epsilon_{ijk} \begin{cases} 0 & \text{any two or three } i=j, j=k, i=k \\ 1 & ijk = 123 \quad 231 \quad 312 \\ -1 & ijk = 132 \quad 321 \quad 122 \end{cases}$$

$$\epsilon_{ijk} \epsilon_{imn} = \delta_{jm} \delta_{kn} - \delta_{jn} \delta_{km}$$

$$\vec{a} \times (\vec{b} \times \vec{c}) = \vec{b}(\vec{a} \cdot \vec{c}) - \vec{c}(\vec{a} \cdot \vec{b})$$

$$(\vec{b} \times \vec{c})_i = \epsilon_{ijk} b_j c_k$$

$$\begin{aligned}(\vec{a} \times (\vec{b} \times \vec{c}))_i &= \epsilon_{ijk} a_j \cdot (\vec{b} \times \vec{c})_k = \epsilon_{ijk} a_j (\epsilon_{kmn} b_m c_n) \\ &= \epsilon_{kij} a_j \epsilon_{kmn} b_m c_n \\ &= \epsilon_{kij} \epsilon_{kmn} a_j b_m c_n \\ &= (\delta_{im} \delta_{jn} - \delta_{in} \delta_{jm}) a_j b_m c_n \\ &= b_i \vec{a} \cdot \vec{c} - c_i \vec{a} \cdot \vec{b}\end{aligned}$$

$$[H, L_x] = [H, L_y] = [H, L_z] = 0$$

$$[L_x, L_y] = i\hbar L_z$$

Can only choose one L_i to find energy eigenstates for