

Parity operator
 $x \rightarrow -x$
 $p \rightarrow -p$

$$[H, P] = 0$$

2 sol^{ns} odd & even

even:
$$\phi^e(x) = \begin{cases} C_1 e^{fx} & x < -a/2 \\ A \cos(qx) & -a/2 < x < a/2 \\ C_1 e^{-fx} & x > a/2 \end{cases}$$

odd:
$$\phi^o(x) = \begin{cases} C_1 e^{fx} & x < -a/2 \\ B \sin(qx) & -a/2 < x < a/2 \\ -C_1 e^{-fx} & x > a/2 \end{cases}$$

Match even sol^{ns}

values: $C_1 e^{-fa/2} = A \cos(qa/2)$

derivatives: $f C_1 e^{-fa/2} = q A \sin(qa/2)$ divide by top

$$f = q \tan(qa/2)$$

transcendental eqⁿ

$$\frac{f(E)}{q(E)} = \frac{\sqrt{-\frac{2mE}{\hbar^2}}}{\sqrt{\frac{2m(E+V_0)}{\hbar^2}}}$$

do some graphs!

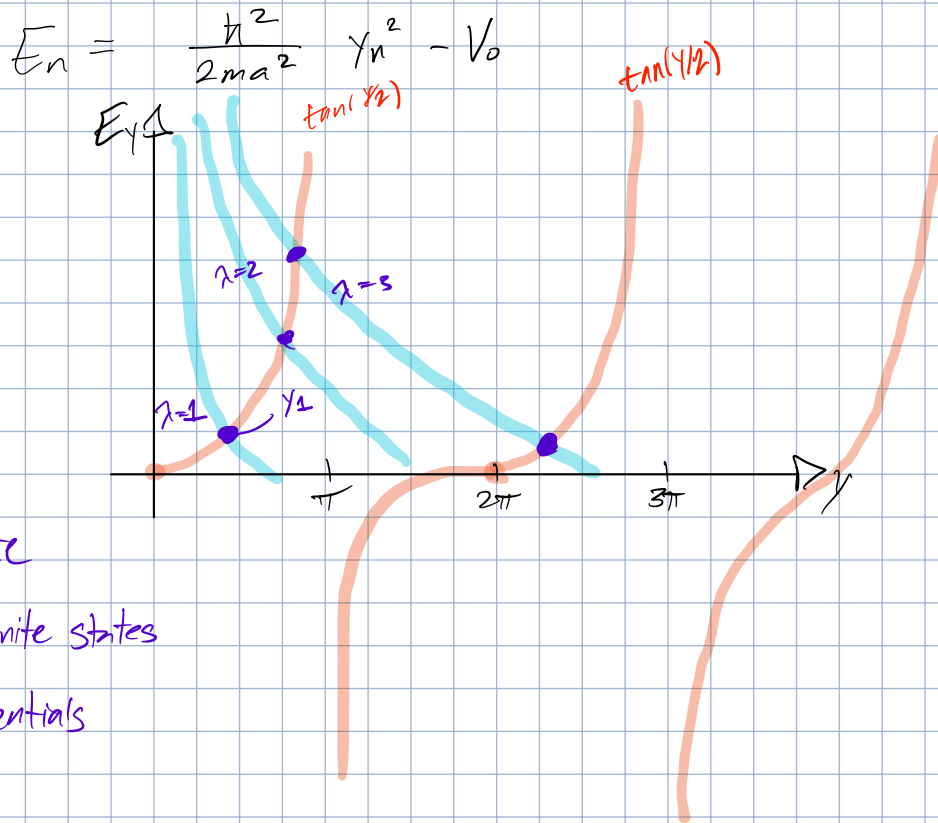
$$\lambda = qa \quad \lambda - \gamma^2 = \frac{2ma^2V_0}{\hbar^2} - \frac{\gamma^2 a^2}{a^2} = -\frac{2mE}{\hbar^2} a^2 = f^2 a^2$$

Strength of how attractive potential is $\leftarrow a \rightarrow$ \leftarrow width \rightarrow V_0 strength

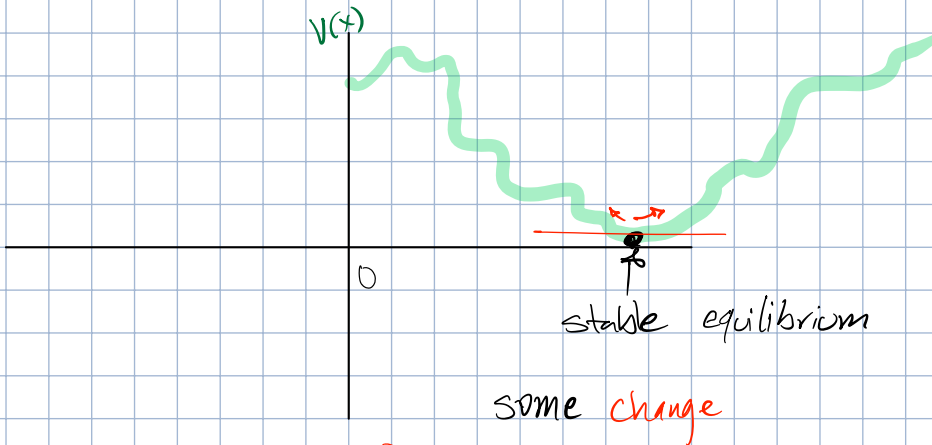
$$\frac{\sqrt{\lambda - \gamma^2}}{\gamma} = \tan(\gamma/2)$$

lowest energy: $E = -V_0 \rightarrow \gamma = 0$
 $E \rightarrow \infty \rightarrow \gamma \rightarrow \infty$

look @ γ solⁿs of
 can get some E 's



Simple Harmonic Oscillator



$$V(x) = \cancel{V(x_{\min})} + \frac{dV}{dx} \Big|_{x_{\min}} (x - x_{\min}) + \frac{1}{2} \frac{d^2V}{dx^2} \Big|_{x_{\min}} (x - x_{\min})^2 + \dots$$

generally nonzero

choose x st. $x_{\min} = 0$

$v(x)$

$$\rightarrow \hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2} m \omega^2 \hat{x}^2$$

$$\frac{1}{2} m \omega^2 \hat{x}^2 = \frac{1}{2} \left. \frac{dV}{dx} \right|_{x_{\min}}$$

Solve 2 ways, useful elsewhere

H-atom
angular momentum

Study relation between classical & Quantum mechanics

Solve time indep. SE
$$-\frac{\hbar^2}{2m} \frac{d^2 \phi_E}{dx^2} + \frac{1}{2} m \omega^2 x^2 \phi_E(x) = E \phi_E(x)$$

find allowed E for normalizable $\phi_E(x)$

① Simplify eqⁿ w/ dimensionless variables

② use asymptotic analysis to find behaviors of solⁿs @ $x \rightarrow \pm \infty$

③ use ② to find series solⁿ to differential eqⁿ

Frobenius method

①

Mass Length Time

$$\left. \begin{aligned} [\omega] &= T^{-1} \\ [m] &= M \\ [\hbar] &= L^2 M T^{-1} \end{aligned} \right\} \left[\frac{m \omega}{\hbar} \right] = L^{-2}$$

can introduce
$$y = \left(\sqrt{\frac{m \omega}{\hbar}} \right) x$$

$$E = \frac{2E}{\hbar \omega} \quad \text{dimensionless energy}$$

$$\frac{d^2 \phi_E}{dy^2} - y^2 \phi_E(y) = -E \phi_E(y)$$

② $f(x) \sim g(x)$ as $x \rightarrow \infty$ if $\lim_{x \rightarrow \infty} \frac{f(x)}{g(x)} = 1$

exe. $f(x) = (1+x)e^x$
 $g(x) = xe^x$ } $f \sim g$ b/c $\lim_{x \rightarrow \infty} \left(\frac{(1+x)e^x}{xe^x} \right) = \lim_{x \rightarrow \infty} \left(\frac{1+x}{x} \right) = \lim_{x \rightarrow \infty} \left(1 + \frac{1}{x} \right) = 1$

math - equivalence relation

set $S = \{a, b, c\}$

\sim is equivalence relation if $a \sim a$
 $a \sim b$ iff $b \sim a$
 $a \sim b, b \sim c \implies a \sim c$

example 1

Set = \mathbb{Z} integers

$n \sim m$ if $n - m = 0 \pmod{2}$

$2 \sim 8$
 $-7 \sim 5$

splits into 2

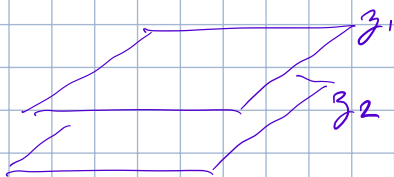
evens = $\{ \dots, -2, 0, 2, 4, \dots \}$
 odds = $\{ \dots, -3, -1, 1, 3, 5, \dots \}$

these are equivalence relation

example 2

Set = \mathbb{R}^3

$(x, y, z) \sim (x', y', z')$ if $z = z'$



divide into classes of \mathbb{R}^3 s based on behavior @ $x \rightarrow \infty$

\rightarrow only lives in equivalence classes of \mathbb{R}^3 s

want simple

$$\phi_E^{\text{asym}}(x) \sim \phi_E(x)$$

↑ solves DE

$$\frac{d^2 \phi_E}{dy^2} - y^2 \phi_E(y) = -\epsilon \phi_E(y)$$

↑
much larger than as $y \rightarrow \pm \infty$

as $y \rightarrow \pm \infty$

$$\frac{d^2 \phi_E^{\text{asym}}}{dy^2} \sim y^2 \phi_E^{\text{asym}}(y)$$

only normalizable so that $x \rightarrow \pm \infty$ doesn't blow up
bc $\int |\phi(x)|^2 = 1$

$$\phi_E^{\text{asym}} = A e^{-y^2/2} \quad \text{for} \quad \frac{d \phi_E^{\text{asym}}}{dy} = -y \phi_E^{\text{asym}}$$

$$\frac{d \phi_E^{\text{asym}}}{dy} = -y A e^{-y^2/2}$$

$$\frac{d \phi_E^{\text{asym}}}{dy} = y^2 A e^{-y^2/2} - A e^{-y^2/2} \sim y^2 A e^{-y^2/2}$$