

an important point about $V \otimes W$

↑
vector space of all linear combinations $v \otimes w$ $v \in V, w \in W$

but not all elements of $V \otimes W$ can be written as $v \otimes w$

not all linear combinations are $v \otimes w$

$$V = \mathbb{C}^2$$

$V \otimes V =$ vector space of all 2×2 complex matrices

choose basis $\begin{pmatrix} 0 & 1 \\ 1 & 1 \end{pmatrix} \in V \otimes V$

but it is not $v \otimes v'$ for any v, v'

$$v \otimes v' = (v_1 \ v_2) \begin{pmatrix} v_1' \\ v_2' \end{pmatrix} = \begin{pmatrix} v_1 v_1' & v_1 v_2' \\ v_2 v_1' & v_2 v_2' \end{pmatrix} \stackrel{?}{=} \begin{pmatrix} 0 & 1 \\ 1 & 1 \end{pmatrix}$$

$$v_1, v_1' = 0$$

v_1 or v_1' is zero

$$v_1 \cdot v_2' = 1$$

$$\text{or } v_2 v_1' = 1$$

↳ but if $v_1 = 0 \rightarrow v_1 v_2' = 0$, but it's 1

contradiction!

↳ but if $v_1' = 0 \rightarrow v_2 v_1' = 0$, but it's 1!

$$V = \mathbb{C}^2$$

these kinds of vectors correspond to entangled quantum state

what about linear operators $\sigma : V \otimes W \rightarrow V \otimes W$

$$O_v : V \rightarrow V$$

$$O_w : W \rightarrow W$$

operators take vector in $V \otimes W$

$$\text{define } (O_v \otimes O_w) : V \otimes W \rightarrow (O_v v) \otimes (O_w w)$$

$$V = \mathbb{C}^2$$

Two photon states

right/left handed circular polarized

Recall from E&M:

$$\vec{E}_{\text{RHCP/LHCP}} = E_0 \text{Re} \left[\left(\hat{x} + i\hat{y} \right) e^{i(kz - \omega t)} \right] \quad i = e^{i\pi/2}$$

$$= E_0 \left(\hat{x} \cos(kz - \omega t) + \hat{y} \cos(kz - \omega t + \pi/2) \right)$$

QM: these correspond to \downarrow -proton states

$$|R\rangle = \frac{1}{\sqrt{2}} (|\vec{x}\rangle + i|\vec{y}\rangle)$$

$$|L\rangle = \frac{1}{\sqrt{2}} (|\vec{x}\rangle - i|\vec{y}\rangle)$$

$$2L|L\rangle = \langle R|R\rangle = 1$$

$$2L|R\rangle = \langle R|L\rangle = 0$$

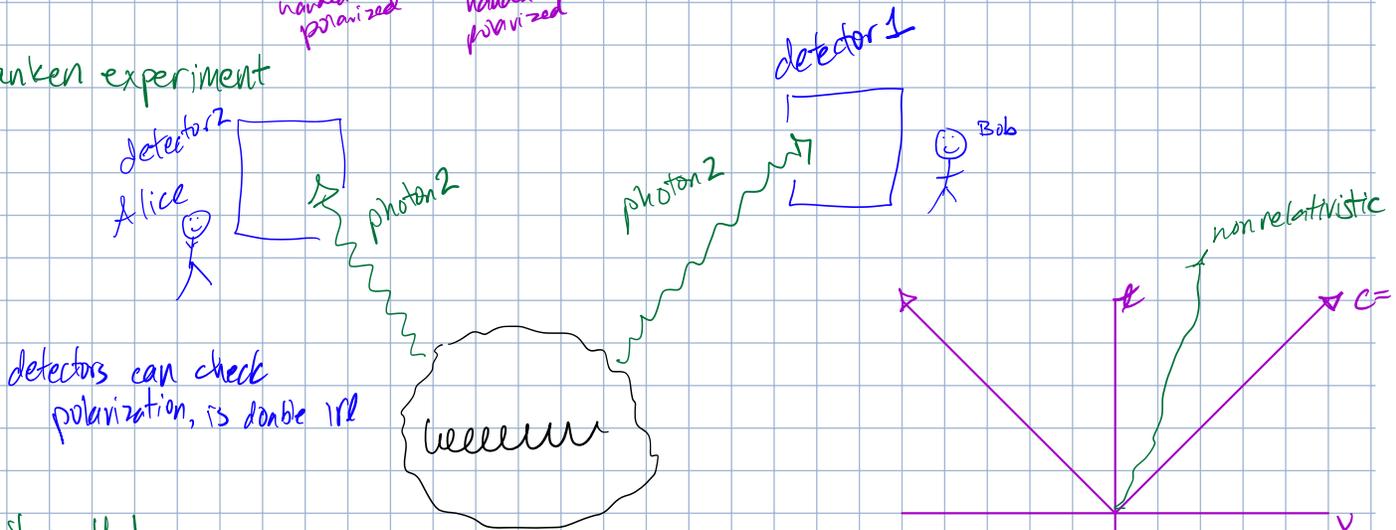
going to look @ state in $\mathbb{C}^2 \otimes \mathbb{C}^2$ (space of 2 photon QM)

$$|EPR\rangle = \frac{1}{\sqrt{2}} (|R\rangle_1 \otimes |L\rangle_2 + |L\rangle_1 \otimes |R\rangle_2)$$

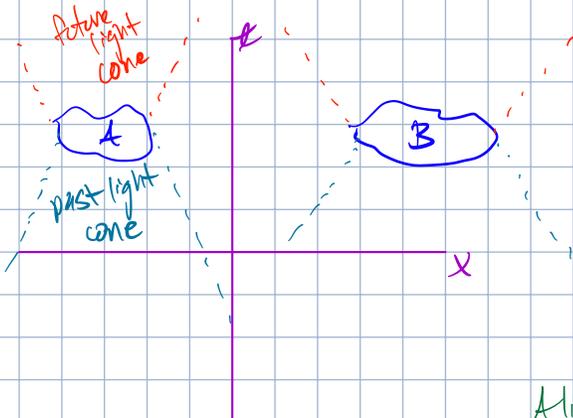
1st photon
2nd photon

1st right handed polarized
2nd left handed polarized
left
right

gedanken experiment



want it so that when Alice measures, has no effect on Bob's measurement
 " Bob Alice



Bob & Alice are spacelike separated
 \rightarrow they lie past each other's future light cone
 so can't interfere w/ each other

$$\left. \begin{matrix} t_A > t_B \\ t_B > t_A \end{matrix} \right\} \text{observer dependent}$$

Alice & Bob can't influence w/ each other while taking measurement

chose x-pol: $\hat{x} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ ± 1 eigenvalue
 chose y-pol: $\hat{y} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ ± 1 eigenvalue $\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

$\sigma_3 |x\rangle = +1 |x\rangle$
 $\sigma_3 |y\rangle = -1 |y\rangle$

measuring ± 1 eigenvalues shows eigenstates being polarized along a direction

RHCP $|R\rangle = \frac{1}{\sqrt{2}} (|x\rangle + i|y\rangle) = \begin{pmatrix} 1 \\ i \end{pmatrix}$
 LHCP $|L\rangle = \frac{1}{\sqrt{2}} (|x\rangle - i|y\rangle) = \begin{pmatrix} 1 \\ -i \end{pmatrix}$

$\sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$

$\sigma_2 - \lambda I = \begin{pmatrix} -\lambda & -i \\ i & -\lambda \end{pmatrix} = \lambda^2 - 1 \rightarrow \lambda = \pm 1$
 $\lambda = 1: \begin{pmatrix} -1 & -i \\ i & -1 \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix} = \begin{pmatrix} -u - iv = 0 \\ iu - v = 0 \end{pmatrix} \rightarrow u = -iv, i(-iv) - v = 0 \rightarrow v = 0, u = 0$
 $\lambda = -1: \begin{pmatrix} 1 & -i \\ i & 1 \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix} = \begin{pmatrix} u - iv = 0 \\ iu + v = 0 \end{pmatrix} \rightarrow u = iv, i(iv) + v = 0 \rightarrow v = 0, u = 0$

$\sigma_2 |R\rangle = +1 |R\rangle$
 $\sigma_2 |L\rangle = -1 |L\rangle$

$\pm 45^\circ$ linear pol. $|\pm\rangle = \frac{1}{\sqrt{2}} (|x\rangle \pm |y\rangle)$

$\sigma_1 =$
 $\sigma_1 |+\rangle = +1 |+\rangle$
 $\sigma_1 |-\rangle = -1 |-\rangle$

Alice & Bob can't influence

\rightarrow assume Alice & Bob can measure $\sigma_i \otimes \sigma_j$ for any i, j

$\sigma_2 \otimes \sigma_3$ measures R,L for photon 1
 \hat{x}, \hat{y} for photon 2
 Bob Alice

Measure R,L (or $\sigma_2 \otimes \sigma_3$) \rightarrow 4 possibilities

Bob	Alice	prob
R	R	0
R	L	$\frac{1}{2}$
L	R	$\frac{1}{2}$
L	L	0

$\text{prob}(R,L) = |(\langle R| \otimes \langle L|) |EPR\rangle|^2$

$= |(\langle R| \otimes \langle L|) \otimes \frac{1}{\sqrt{2}} (|R\rangle_2 \otimes |L\rangle_2 + |L\rangle_2 \otimes |R\rangle_2) |^2$

$= |(\langle R| \otimes \langle L|) \otimes \frac{1}{\sqrt{2}} (|R\rangle_2 \otimes |L\rangle_2) + (\langle R| \otimes \langle L|) \otimes \frac{1}{\sqrt{2}} (|L\rangle_2 \otimes |R\rangle_2) |^2$

$= | \frac{1}{\sqrt{2}} (\langle R| \otimes \langle L|) \otimes (|R\rangle_2 \otimes |L\rangle_2) + \frac{1}{\sqrt{2}} (\langle R| \otimes \langle L|) \otimes (|L\rangle_2 \otimes |R\rangle_2) |^2$

$= | \frac{1}{\sqrt{2}} (\langle R|R\rangle_2 \otimes \langle L>L\rangle_2) + \frac{1}{\sqrt{2}} (\langle R>L\rangle_2 \otimes \langle L>R\rangle_2) |^2$

$= | \frac{1}{\sqrt{2}} (1 \cdot 1) + \frac{1}{\sqrt{2}} (0 \cdot 0) |^2 = | \frac{1}{\sqrt{2}} |^2 = \frac{1}{2}$

$$= \frac{1}{\sqrt{2}} / 2 = \frac{1}{2}$$

Alice & Bob each see a random distribution: $\frac{1}{2} R$ $\frac{1}{2} L$

if Bob gets R, Alice gets L

despite being space indpt.

Einstein: spukhafte Fernwirkung
spooky long dist. effect

EPR \rightarrow entangled states