

time indep SE

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \phi(x) + U(x) \phi(x) = E \phi(x)$$

constant or piecewise constant

$$V(x) = U_0$$

$$\frac{d^2 \phi}{dx^2} = -\frac{2m}{\hbar^2} (V_0 - E) \phi(x)$$

Solutions: $Ae^{ikx} + Be^{-ikx}$

if $\frac{2m}{\hbar^2} (V_0 - E) < 0$

→ particle where $E > V_0$

→ $k^2 = \frac{2m}{\hbar^2} (V_0 - E)$

or

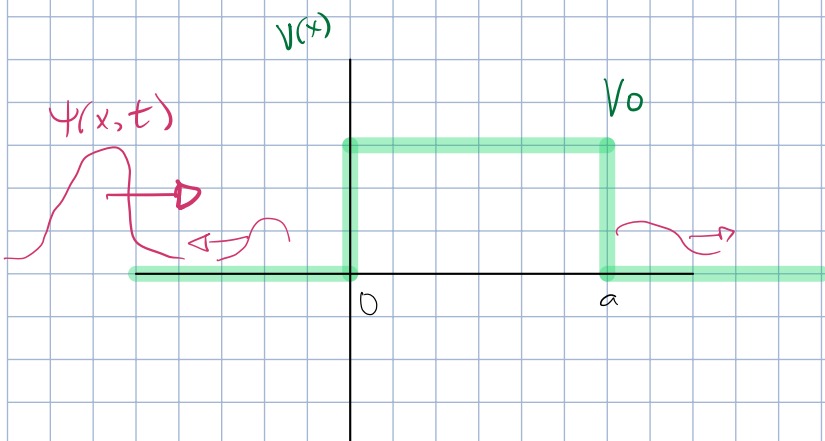
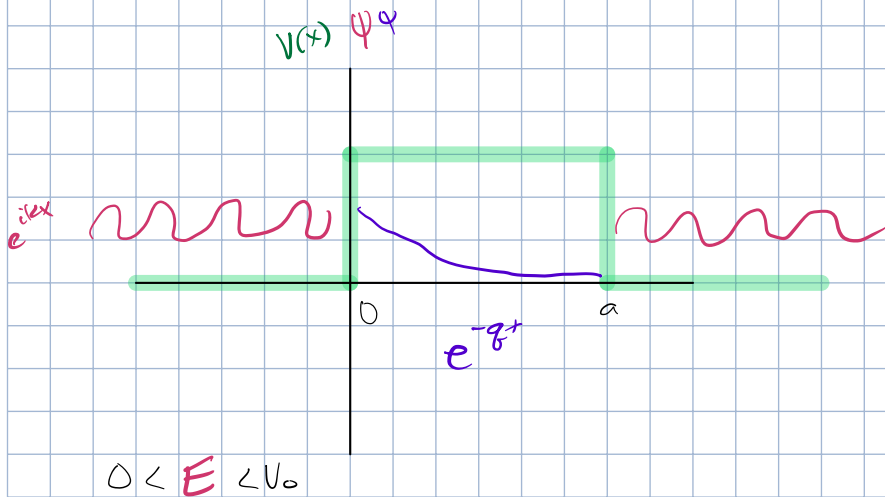
$Ce^{iqx} + De^{-iqx}$

if $\frac{2m}{\hbar^2} (V_0 - E) > 0$

→ " " $E < V_0$

not allowed for $x \rightarrow +\infty$

→ $q^2 = \frac{2m}{\hbar^2} (V_0 - E)$



solving $\psi(x,t)$ is hellish hard

sometimes reflected or transmitted

Probability Density

$$P(x,t) = |\psi(x,t)|^2 = |\langle x | \psi(t) \rangle|^2$$

around a region $\int P(x,t) dx$

norm. $\int_{-\infty}^{\infty} P(x,t) dx = 1$

but locally, probability can change

$\leadsto \frac{\partial P}{\partial t} ?$

$$\begin{aligned} \rightarrow \frac{\partial P(x,t)}{\partial t} &= \frac{\partial}{\partial t} \psi^*(x,t) \psi(x,t) \\ &= \frac{\partial \psi^*}{\partial t} \psi + \psi^* \frac{\partial \psi}{\partial t} \end{aligned}$$

we also know time dependent SE

$$\rightarrow \frac{\partial \psi}{\partial t} = \frac{1}{i\hbar} \left(-\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + U(x) \psi \right)$$

$$\frac{\partial \psi^*}{\partial t} = \frac{-1}{i\hbar} \left(-\frac{\hbar^2}{2m} \frac{\partial^2 \psi^*}{\partial x^2} + U(x) \psi^* \right)$$

\rightarrow assuming $U(x) = U^*(x)$ needed for $H = H^\dagger$

$$\begin{aligned} \frac{\partial P}{\partial t} &= \frac{\hbar}{2im} \left(\psi \frac{\partial^2 \psi^*}{\partial x^2} - \psi^* \frac{\partial^2 \psi}{\partial x^2} \right) \\ &= \frac{-\partial}{\partial x} \left(\frac{\hbar^2}{2mi} \left(\psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} \right) \right) \end{aligned}$$

potential goes away

$$j(x,t) =$$

$$\frac{\partial P}{\partial t} = -\frac{\partial}{\partial x} j$$

$$\frac{\partial P}{\partial t} + \frac{\partial j}{\partial x} = 0 \quad \text{conservation!}$$

like in EM

ρ = charge density
 j = current density

$$\rightarrow \frac{\partial \rho}{\partial t} + \nabla \cdot j = 0$$

how much
is flowing



$$\int_a^b P(x,t) dx \rightarrow \text{prob of it being there}$$

$$\frac{\partial}{\partial t} \int_a^b P(x,t) dx = - \int_a^b \frac{\partial j}{\partial x} dx$$
$$= j(a) - j(b)$$

$$\psi_p(x) = e^{\frac{ipx}{\hbar}} \quad \text{momentum eigenstate}$$

assume no $U(x)$

what is prob. current of $\psi_p(x)$?

$$j = \frac{\hbar}{2im} \left(\psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} \right)$$

$$= \frac{\hbar}{2im} \left(e^{\frac{-ipx}{\hbar}} \cdot \frac{ip}{\hbar} \cdot e^{\frac{ipx}{\hbar}} - \dots \right)$$

$$= \frac{\hbar}{2im} \left(\psi^* \frac{\partial \psi}{\partial x} - \left(\psi^* \frac{\partial \psi}{\partial x} \right)^* \right)$$

$$= \frac{\hbar}{2im} \left(e^{\frac{-ipx}{\hbar}} \cdot \frac{ip}{\hbar} \cdot e^{\frac{ipx}{\hbar}} + e^{\frac{ipx}{\hbar}} \cdot \frac{-ip}{\hbar} \cdot e^{\frac{-ipx}{\hbar}} \right)$$

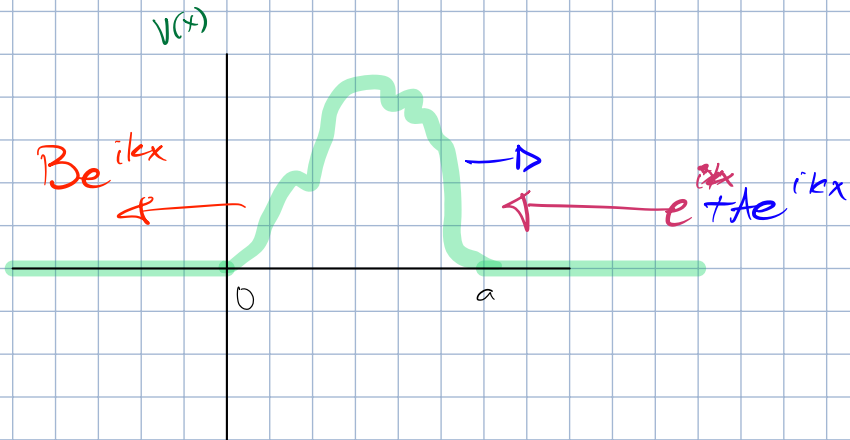
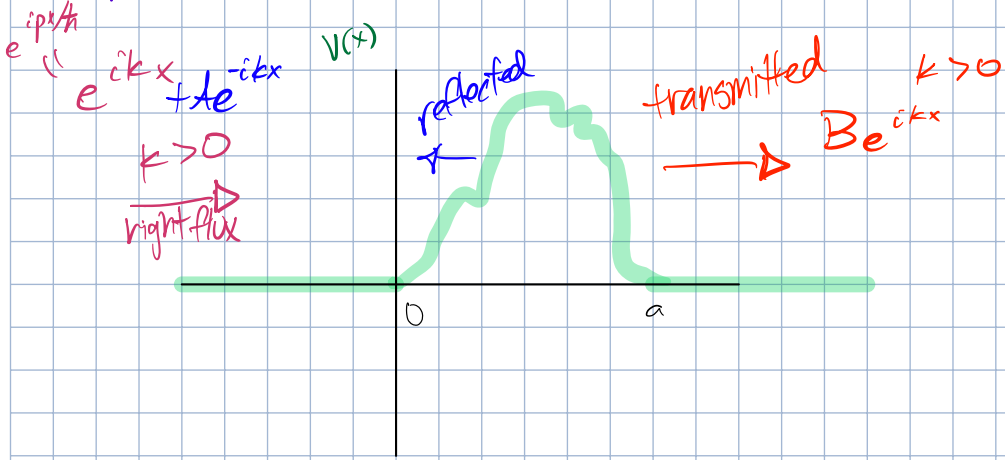
$$= \frac{\hbar}{2im} \left(\frac{ip}{\hbar} + \frac{ip}{\hbar} \right) = \frac{p}{m}$$

$p > 0 \quad j > 0$ flux to the right

prob. density - amt water in river

particles moving to the right

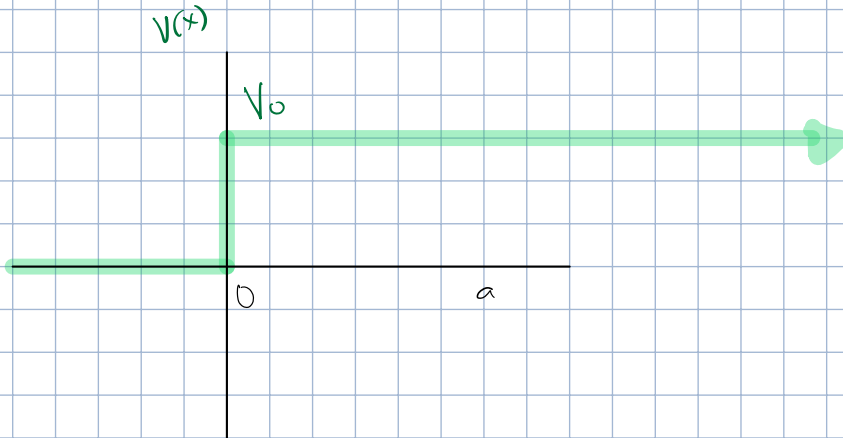
$$\Psi_{\text{wavepacket}}(x,t) = \int e^{ipx} \left(\pi - \frac{i\hbar^2 t}{2m\hbar} \right) \phi(p) dp$$



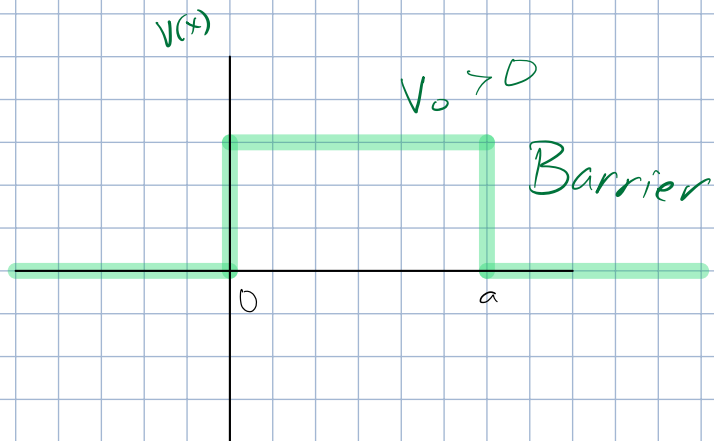
A & B hold physics

Cool examples

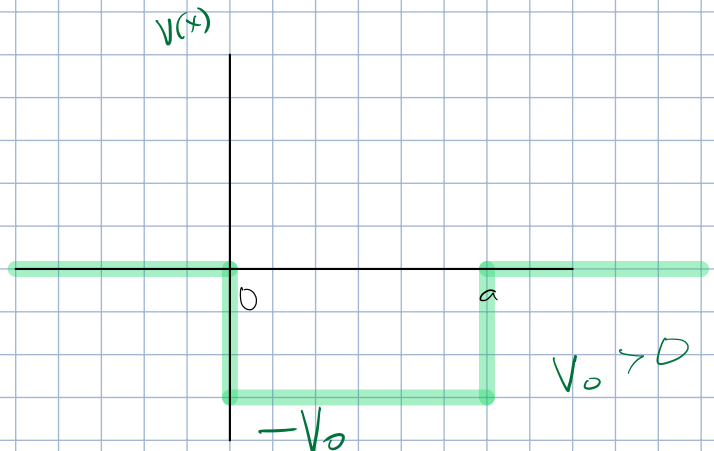
① Step potential



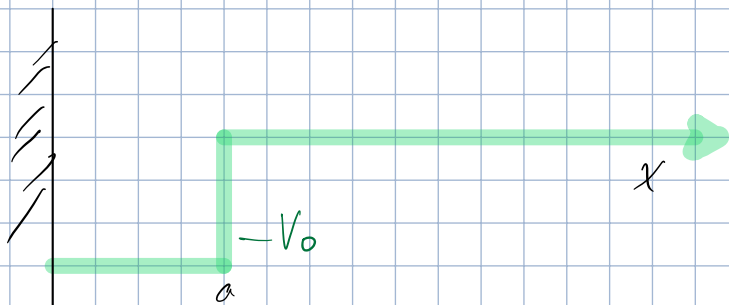
② Repulsive Square Well



③ Attractive Square Well



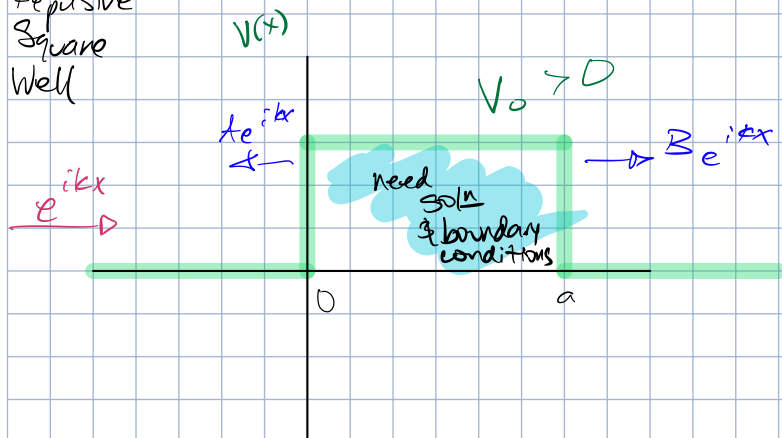
④



spherical coords

Part 2

Repulsive Square Well



No time dependence in $\Psi(x,t)$ (turns out to just change phase)

$$\text{so } \frac{\partial \Psi}{\partial t} = 0$$

$$\rightarrow \frac{\partial \Psi}{\partial x} = 0$$

for $x < 0$ $\Psi = e^{ikx} + Ae^{-ikx}$

$$j = \frac{\hbar k}{2m} \left(\Psi^* \frac{\partial \Psi}{\partial x} - \Psi \frac{\partial \Psi^*}{\partial x} \right)$$

$$j = \frac{\hbar k}{m} (1 - |A|^2)$$

$$\Psi^* \frac{\partial \Psi}{\partial x} = (e^{-ikx} + A^* e^{ikx}) \cdot (ik e^{ikx} - ik e^{-ikx})$$

$$= ik (1 - |A|^2 + A^* e^{2ikx} - A e^{-2ikx})$$

\downarrow real \downarrow complex
 with i , pure imaginary pure real

for $x > a$ $\Psi = B e^{ikx}$

$$j = \frac{\hbar k}{m} |B|^2 \quad \text{nothing goes other way}$$

but we subtract by complex conjugate

so only real stays

$$1 - |A|^2 = |B|^2$$

in - reflected = + transmitted, conservation

convention: $|B|^2 = T$ = transmission probability

$|A|^2 = R$ = reflection probability

$$1 = R + T$$

for SE, need solⁿ in all regions

$$k^2 = \frac{2m(V_0 - E)}{\hbar^2}$$

$$q^2 = -\frac{2m(V_0 - E)}{\hbar^2} = \frac{2m(E - V_0)}{\hbar^2}$$

I. $x < 0 \rightarrow \Psi_I = e^{ikx} + A e^{-ikx}$

need to decide if $E > V_0$ (scattering over barrier) . would classically slow down
 assume $E > V_0$

II. $0 < x < a \rightarrow \Psi_{II} = C e^{iqx} + D e^{-iqx}$

III. $x > a \rightarrow \Psi_{III} = B e^{ikx}$

2 conditions

Demand $\Psi(x,t)$ is continuous in x

potentials are idealized

at $x = 0, a$

$$\Psi_I(0) = \Psi_{II}(0)$$

$$\Psi_{II}(a) = \Psi_{III}(a)$$

$$\lim_{\epsilon \rightarrow 0} \left(\frac{d\psi}{dx} \Big|_{\epsilon^-} - \frac{d\psi}{dx} \Big|_{\epsilon^+} \right) = \lim_{\epsilon \rightarrow 0} \left(\int_{-\epsilon}^{\epsilon} \frac{d^2\psi}{dx^2} dx \right) = \lim_{\epsilon \rightarrow 0} \left(\int_{-\epsilon}^{\epsilon} \frac{2m}{\hbar^2} (V(x) - E) dx \right) = 0$$

$$\frac{d\psi_{I(0)}}{dx} = \frac{d\psi_{II(0)}}{dx}$$

not singular, so we chillin

$$\frac{d\psi_{II(a)}}{dx} = \frac{d\psi_{III(a)}}{dx}$$

Now algebra & calculus

$$x=0 \quad \psi \quad 1+A=C+D$$

$$\psi' \quad ik(1-A) = iq(C-D)$$

$$x=a \quad \psi \quad Ce^{iqa} + De^{-iqa} = Be^{ika}$$

$$\psi' \quad iq(Ce^{iqa} - De^{-iqa}) = ikBe^{ika}$$

$$R = |A|^2 = \frac{(k^2 - q^2)^2 \sin^2(qa)}{4k^2q^2 + (k^2 - q^2)^2 \sin^2(qa)}$$

$$T = |B|^2 = \frac{4k^2q^2}{4k^2q^2 + (k^2 - q^2)^2 \sin^2(qa)}$$

discussion

as $V_0 \rightarrow 0$, $q^2 \rightarrow k^2$, $R \rightarrow 0$, $T \rightarrow 1$

$$T = \frac{1}{1 + \frac{(k^2 - q^2)^2}{4k^2q^2} \sin^2(qa)}$$

T varying w/E

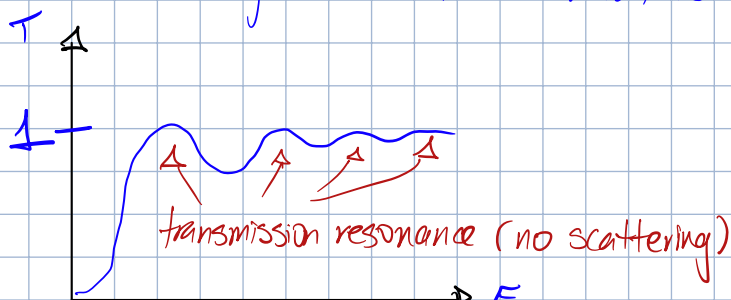
$$q = \sqrt{2m(V_0 - E)/\hbar^2}$$

\sin^2 varies between 0 & 1

when $\sin^2 qa = 0 \rightarrow qa = 0, \pi, 2\pi, \dots$

$\rightarrow T = 1$

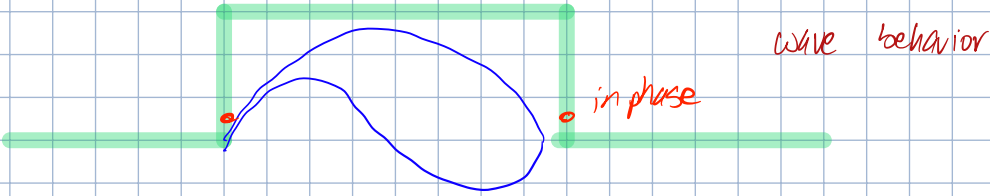
certain energies where all of wave is transmitted



looking e $0 < x < a$ $e^{i\pi x}$

$$\lambda = \frac{2a}{n}$$

$$qa = n\pi \rightarrow \lambda = \frac{2a}{n}$$



Ramsauer-Townsend effect

low energy scattering electrons of noble gases

closed shells of electrons



Spherical potential