

Ehrenfest's thm
 derive uncertainty principle
 1-d potential

free particle $U(x)$, SHO $V(x) = \frac{1}{2}kx^2 = \frac{1}{2}m\omega^2x^2$, finite square well $\begin{matrix} a & -2a \\ | & - & | \\ \hline & & \end{matrix}$

Ehrenfest's Thm

$$\frac{dx}{dt} = \{x, H\} \quad \frac{dp}{dt} = \{p, H\} \quad \text{Poisson brackets}$$

↙ Hermitian

$$\langle \psi | \hat{O} | \psi \rangle = \text{expectation value} = \sum \lambda_i p_i$$

$|\psi(x,t)\rangle$ obeying $\hat{H} \psi(x,t) = i\hbar \frac{\partial \psi}{\partial t}$

$$\langle x | \psi(t) \rangle = \psi(x,t)$$

$$\langle \psi(t) | \hat{x} | \psi(t) \rangle = \int \psi^*(x,t) \cdot x \cdot \psi(x,t) dx$$

↑
 limits depend on problem
 assume it vanishes @ limits

∫ other finite potentials

want to compute its time derivative

$$\frac{d}{dt} \langle \psi(t) | \hat{x} | \psi(t) \rangle = \int \left(\frac{\partial \psi^*}{\partial t} \hat{x} \psi + \psi^* \hat{x} \frac{\partial \psi}{\partial t} \right) dx$$

$$\frac{\partial \psi}{\partial t} = \frac{i}{\hbar} \hat{H} \psi$$

$$\frac{\partial \psi^*}{\partial t} = \frac{i}{\hbar} (\hat{H} \psi)^*$$

$\hat{H} = \hat{H}^*$ b/c
 † b/c integration by parts

$$= \int \left(\frac{i}{\hbar} (\hat{H} \psi)^* \hat{x} \psi - \frac{i}{\hbar} \psi^* \hat{x} \hat{H} \psi \right) dx$$

$$= \int \frac{i}{\hbar} \psi^* \hat{H} \hat{x} \psi - \frac{i}{\hbar} \psi^* \hat{x} \hat{H} \psi dx$$

$$= \int \psi^* \left(\frac{i}{\hbar} \hat{H} \hat{x} - \frac{i}{\hbar} \hat{x} \hat{H} \right) \psi dx$$

$$= \int \psi^* \left(-\frac{i}{\hbar} (\hat{x} \hat{H} - \hat{H} \hat{x}) \right) \psi dx$$

$$= \int \psi^*(x,t) \left(-\frac{i}{\hbar} [\hat{x}, \hat{H}] \right) \psi(x) dx$$

$$= \langle \psi(t) | -\frac{i}{\hbar} [\hat{x}, \hat{H}] | \psi(t) \rangle$$

$$\frac{d\langle x \rangle}{dt} = \frac{i}{\hbar} \langle [\hat{x}, \hat{H}] \rangle$$

$$\{f, g\} \rightarrow -\frac{i}{\hbar} [\hat{f}, \hat{g}] \quad \dagger \text{ avg. } \langle \psi | \quad | \psi \rangle$$

$$\frac{d\langle p \rangle}{dt} = -\frac{i}{\hbar} \langle [\hat{p}, \hat{H}] \rangle$$

→ with $\hat{p} \rightarrow \hat{p} | \psi(t) \rangle$

$$\frac{d}{dt} \langle \psi(t) | \hat{p} | \psi(t) \rangle = \frac{i}{\hbar} \langle \psi(t) | [\hat{p}, \hat{H}] | \psi(t) \rangle$$

suppose $\hat{H} = \frac{\hat{p}^2}{2m} + V(x)$

Check
this math

$$[\hat{x}, \hat{H}] = \hat{x}\hat{H} - \hat{H}\hat{x} = \frac{1}{2m} [\hat{x}, \hat{p}^2] = \frac{1}{2m} (\hat{x}\hat{p}\hat{p} - \hat{p}\hat{p}\hat{x})$$

$$[\hat{x}, \hat{p}] = i\hbar \quad \hat{x}\hat{p} - \hat{p}\hat{x} = i\hbar \quad \text{or} \quad \hat{x}\hat{p} = \hat{p}\hat{x} + i\hbar$$

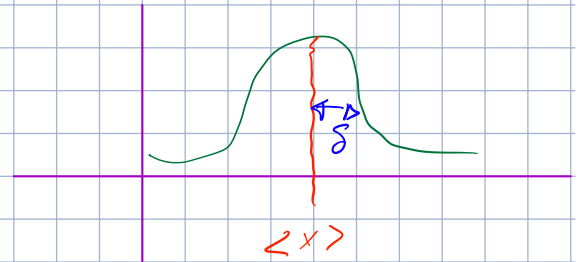
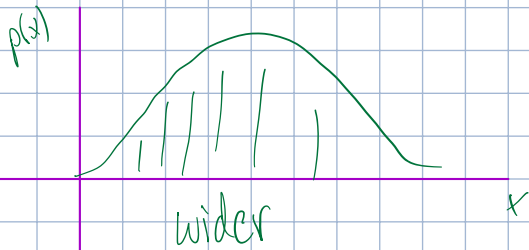
$$\hat{x}\hat{p}\hat{p} = (\hat{p}\hat{x} + i\hbar)\hat{p} = \hat{p}\hat{p}\hat{x} + i\hbar\hat{p} + i\hbar\hat{p} = \hat{p}\hat{p}\hat{x} + 2i\hbar\hat{p}$$

$$[\hat{x}, \hat{H}] = \frac{1}{2m} (2i\hbar\hat{p}) = -\frac{i\hbar}{m} \hat{p} \quad \text{so} \quad \frac{d}{dt} \langle \psi(t) | \hat{x} | \psi(t) \rangle = \frac{1}{m} \langle \psi(t) | \hat{p} | \psi(t) \rangle$$

$p=mv$ in avg sense

$$\text{check } [\hat{p}, \hat{H}] = -i\hbar \frac{dV}{dx} \quad \frac{d}{dt} \langle \psi(t) | \hat{p} | \psi(t) \rangle = \langle \psi(t) | -\frac{dV}{dx} | \psi(t) \rangle$$

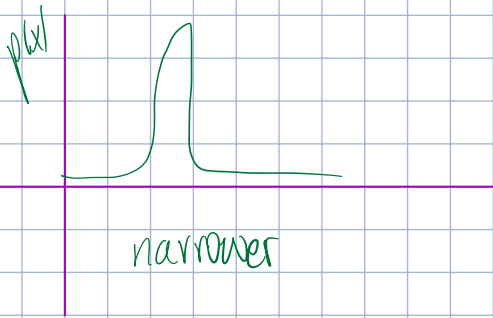
$$[\hat{x}, \hat{p}] = -i\hbar$$



want to see $x - \delta$

or $x + \delta$

but this would cancel out on either side, so we square it



just some #

if A is what you measure some probability $\langle A \rangle$ for its avg

dispersion of probability.
= uncertainty

$$(\Delta A)^2 = \langle (A - \langle A \rangle)^2 \rangle$$

$$= \langle A^2 - 2\langle A \rangle A + \langle A \rangle^2 \rangle$$

$$= \langle A^2 \rangle - \langle A \rangle^2$$

\uparrow \uparrow
 avg. prob.

in QM,

$$\langle \psi(t) | \hat{O} | \psi(t) \rangle = \text{avg value of } \hat{O}$$

$$\langle \psi(t) | \hat{O}^2 | \psi(t) \rangle = \text{avg value of } \hat{O}^2$$

\rightarrow avg of x^2

avg of prob. x^2

$$(\Delta x)^2 = \langle \psi(t) | \hat{x}^2 | \psi(t) \rangle - \langle \psi(t) | \hat{x} | \psi(t) \rangle^2$$

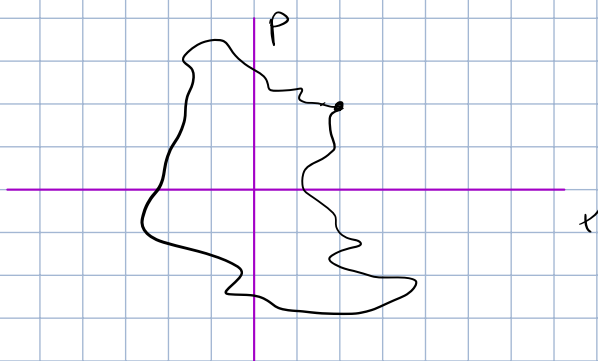
$$(\Delta p)^2 = \langle \psi(t) | \hat{p}^2 | \psi(t) \rangle - \langle \psi(t) | \hat{p} | \psi(t) \rangle^2$$

proof of

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

for any normalized $\psi(x,t)$

classical mechanics: phase space (x,p)



in QM: phase space is chopped up

