

MWE in conductors!

$$\nabla \cdot \tilde{\mathbf{E}} = \frac{\rho_f}{\epsilon}$$

$$\nabla \cdot \tilde{\mathbf{B}} = 0$$

$$\nabla \times \tilde{\mathbf{E}} = -\dot{\tilde{\mathbf{B}}}$$

$$\nabla \times \tilde{\mathbf{B}} = \epsilon \mu \dot{\tilde{\mathbf{E}}} + \mu \tilde{\mathbf{J}}_f$$

$$= \sigma_f \tilde{\mathbf{E}}$$

conduct. w. ity

Wave Eqn

$$\nabla^2 \tilde{\mathbf{E}} = \mu \epsilon \dot{\dot{\tilde{\mathbf{E}}}} + \mu \sigma \dot{\tilde{\mathbf{E}}}$$

$$\nabla^2 \tilde{\mathbf{B}} = \mu \epsilon \dot{\dot{\tilde{\mathbf{B}}}} + \mu \sigma \dot{\tilde{\mathbf{B}}}$$

$$\rho_f = 0$$

$$\nabla \cdot \tilde{\mathbf{J}}_f = -\dot{\rho}_f$$

$$\nabla \cdot (\sigma \tilde{\mathbf{E}}) = -\dot{\rho}_f$$

$$\frac{\sigma}{\epsilon} \rho_f = -\dot{\rho}_f$$

$$\rightarrow \rho_f(t) = \rho_f(t=0) e^{-(\sigma/\epsilon)t}$$

LHS: $\nabla^2 \tilde{\mathbf{E}} = -k^2 \tilde{\mathbf{E}}$

RHS: $\mu \epsilon \dot{\dot{\tilde{\mathbf{E}}}} = -\omega^2 \mu \epsilon \tilde{\mathbf{E}}$
 $\mu \sigma \dot{\tilde{\mathbf{E}}} = -i \gamma \omega \tilde{\mathbf{E}}$

$$\rightarrow -k^2 = -\omega^2 \mu \epsilon - i \mu \sigma \omega$$

say k is complex: $\tilde{k} = \pm (k \pm i \kappa)$

$$k^2 = \frac{\omega^2}{2} \epsilon \mu \left[\sqrt{1 + \left(\frac{\sigma}{\epsilon \omega}\right)^2} + 1 \right]$$

$$\kappa^2 = \frac{\omega^2}{2} \epsilon \mu \left[\sqrt{1 + \left(\frac{\sigma}{\epsilon \omega}\right)^2} - 1 \right]$$

plug in & solve quadratic

$$\left[\frac{\sigma}{\epsilon}\right] \sim \tau^{-1}$$

$$Q^{-1} \equiv \left(\frac{\sigma}{\epsilon \omega}\right) \quad \text{"quality factor"} \quad \text{dimensionless}$$

relates displacement vs conductor current
 $\sigma \dot{\tilde{\mathbf{E}}}$ vs $\sigma \tilde{\mathbf{E}}$

① $Q \gg 1 \rightarrow$ insulator
 conductivity small
 frequency of waves is large

② $Q \ll 1 \rightarrow$ conductor

$$[\text{wave vector}] = L^{-1} = [k] = [\kappa]$$

$$\frac{1}{k^2} = d_{\text{skin}}^2 \quad (\text{length}) = \left(\frac{2}{\omega^2 \epsilon \mu}\right) \cdot \left[\sqrt{1 + \frac{1}{Q^2}} - 1\right]^{-1}$$

$$d_{\text{skin}} = \left(\frac{2}{\epsilon \mu \omega^2}\right)^{1/2} \cdot \left[\sqrt{1 + \frac{1}{Q^2}} - 1\right]^{-1/2}$$

① $Q \gg 1: \quad \sigma \ll \omega \epsilon$

$$k \approx \frac{\sigma}{2} \sqrt{\frac{\mu}{\epsilon}} \rightarrow d_{\text{skin}} \approx \frac{2}{\sigma} \sqrt{\frac{\epsilon}{\mu}}$$

$$\textcircled{2} Q \ll 1: \quad \sigma \gg \omega \epsilon$$

$$k \approx \sqrt{\frac{\mu \omega^2}{2}} \quad \rightarrow \quad \text{damp} \approx \sqrt{\frac{2}{\sigma \mu \omega}}$$

$$\hat{E}(z, t) = \hat{E}_0 \underbrace{e^{-kz}}_{\text{???}} e^{i(kz - \omega t)} \quad \hat{x}$$

$$\hat{B}(z, t) = \frac{k^2}{\omega} \hat{E}_0 \underbrace{e^{-kz}} e^{i(kz - \omega t)} \quad \hat{y}$$

killing \hat{B} & \hat{E} going into material (z grows)
act as damping term

$$\nabla \times \hat{E} = (+ik - k) \hat{E}$$

$$-d_t \hat{B} = +i\omega \hat{B}$$

$$\rightarrow \hat{B}_0 = \frac{1}{\omega} |ik + k| \hat{E}_0$$

just a phase factor! E & B related by complex num
phase no longer fixed!

$$\tilde{K} = K e^{i\phi} = K (\cos \phi + i \sin \phi)$$

$$K^2 = |\tilde{K}|^2 = k^2 + k^2 = \omega^2 \left(\epsilon \mu \sqrt{1 + (1/Q)^2} \right)$$

$$\phi \equiv \tan^{-1} \left(\frac{k}{k} \right)$$

$$\tilde{K} = k + ik$$