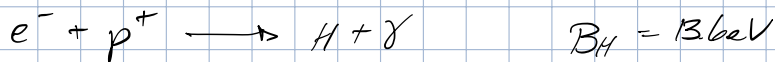


last time: μ -recombination in equilibrium



match μ 's: $\mu_e + \mu_p = \mu_H + \cancel{\mu_\gamma}$

$T \ll m$

take ratios

$$\begin{aligned} \left(\frac{n_H}{n_e n_p}\right)_{\text{eq}} &= \frac{g_H}{g_e g_p} \cdot \left(\frac{m_H}{m_e m_p} \frac{2\pi}{T}\right)^{3/2} \frac{\exp\left[-\frac{m_H c^2}{T}\right]}{\exp\left[-\frac{m_p c^2}{T}\right] \exp\left[-\frac{m_e c^2}{T}\right]} \\ &= \left(\frac{m_H}{m_e m_p}\right)^{3/2} e^{+B_H/T} \\ &= \left(\frac{2\pi}{m_e T}\right)^{3/2} e^{+B_H/T} \end{aligned}$$

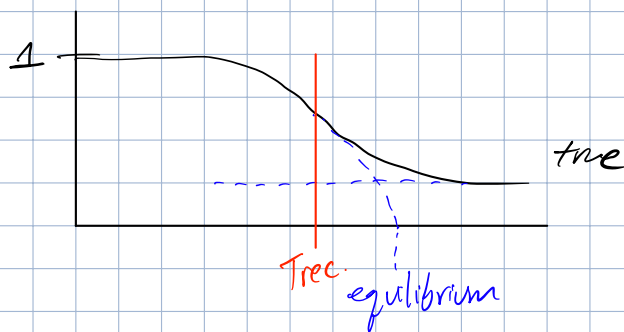
$$B_H \equiv m_p + m_e - m_H = 13.6 \text{ eV}$$

ionized fraction

$$X_e \equiv \frac{n_e}{n_b} = \frac{n_e}{n_p + n_H}$$

(can ignore neutrons $\rightarrow 1/8$)

X_e starts @ 1 & degrades over time



$$T_{\text{recombination}} = 0.3 \text{ eV}$$

charge neutrality: $n_p = n_e$

$$n_H = n_b (1 - X_e)$$

$\hookrightarrow 0 \rightarrow m_e$

$$n_b = \eta \cdot n_\gamma = \eta \frac{2 I(3)}{\pi^2} T^3$$

$$\left(\frac{n_H}{n_e n_p}\right)_{\text{eq}} = \left(\frac{n_H}{n_e}\right)_{\text{eq}}^2$$

$$\rightarrow \frac{1 - X_e}{X_e^2} = \frac{1 - X_e}{(n_e/n_b)^2} = \frac{n_b^2 (1 - X_e)}{n_e^2} = \frac{n_b \cdot n_b (1 - X_e)}{n_e^2}$$

$$= \frac{n_H}{n_e^2} n_b = \frac{n_H}{n_e^2} \cdot \eta \cdot n_\gamma$$

$$\frac{1 - X_e}{X_e^2} = \left(\frac{2\pi}{m_e T}\right)^{3/2} e^{B_H/T} \eta \cdot n_\gamma = \eta \cdot \frac{2 I(3)}{\pi^2} \cdot \left(\frac{2\pi T}{m_e}\right)^{3/2} e^{B_H/T}$$

Timeline ...	γ -dec	$e^+e^- \rightarrow \gamma\gamma$	BBN (He formation)	Recombination (H formation)	γ -decoupling (CMB)
	$T_\gamma \sim \text{MeV}$	$T_\gamma \neq T_e$ $T_\gamma \sim \text{MeV}/2 \text{ MeV}$	$T \sim \text{keV}$	$T \sim 0.3 \text{ eV}$	$T \sim \text{eV}$ (matter domination)

CMB formation

$$e^- + \gamma \leftrightarrow e^- + \gamma$$

photons can't go far... till now!

$$\rho_{\text{tot}} = \underbrace{\rho_\nu + \rho_\gamma + \rho_{\text{DM}}}_{\text{relativistic}} + \underbrace{\rho_B + \rho_e}_{\text{nonrelativistic } T \ll m}$$

$\rho_\nu + \rho_\gamma \ll \rho_{\text{DM}} + \rho_B + \rho_e$

$$\rho_B \gg \rho_e$$

Hubble rate: $H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi}{3} G \rho_{\text{tot}}$

$$= \frac{8\pi}{3} G (m_p \cdot n_b \cdot b)$$

$$= 16\pi G m_p \cdot n \cdot n_\gamma$$

$$= 16\pi G m_p \cdot n \cdot \left(\frac{2J(3)}{\pi^2} T^3\right) = G \frac{32}{\pi} \cdot n \cdot J(3) \cdot m_p T^3 = H^2$$

Rate of reaction: when can it go 1 Hubble time before hitting another e^- ?

$$R_\gamma = n_e \cdot \sigma_T$$

$$= n \cdot \left(\frac{2J(3)}{\pi^2} T^3\right) \cdot \left(\frac{8\pi}{3} \frac{\alpha^2}{m_e^2}\right)$$

$$\sigma_T \equiv \frac{8\pi}{3} \frac{\alpha^2}{m_e^2} \quad \alpha = 1/137$$

(Thomson scattering)

(if all e^- 's are free)
not true tho 1/2 H's are being found

ion-changed particle

$$R_\gamma \implies R_\gamma^{\text{ion}} = X_e \cdot R_\gamma$$

set $H \approx R_\gamma^{\text{ion}}$

$$\left(\frac{32}{\pi} G n J(3) m_p T^3\right)^{1/2} = X_e \left(n \frac{2J(3)}{\pi^2} T^3\right) \cdot \left(\frac{8\pi}{3} \frac{\alpha^2}{m_e^2}\right)$$

$$T_{\text{dec}} \approx 0.27 \text{ eV}$$

about 390,000 years after Big Bang

- ① ok - but dark matter
- ② non equilibrium physics \star

Boltzman Equation

$$f_{\text{eq}}(p, t) = \frac{g}{\exp[-\beta(E_p - \mu)] \pm 1} \quad \text{max entropy}$$

$$N = \int dx^3 \frac{d^3p}{(2\pi)^3} f(\vec{p}, \vec{x}, t) \quad \text{isotropic \& homogeneous Universe}$$

$$n = \int \frac{d^3p}{(2\pi)^3} f(p, x, t) \quad f - \text{phase space distribution}$$

left hand side

$$\frac{d}{dt} [f(p, t)] = \frac{df}{dt} + \frac{dp}{dt} \cdot \frac{df}{dp}$$

$$= \frac{\partial f}{\partial t} - H \cdot p(t) \cdot \frac{\partial f}{\partial p}$$

$$p(t) = \frac{p_0}{a(t)} \quad \left\{ \begin{array}{l} \text{why inverse 1?} \\ \text{why not } \frac{p_0}{a^2} \end{array} \right.$$

$$\frac{\partial p}{\partial t} = -\frac{\dot{a}}{a^2} p_0$$

integrate over $d^3p/(2\pi)^3$

$$\int \frac{d^3p}{(2\pi)^3} \left(\frac{df}{dt} \right) = \int \frac{d^3p}{(2\pi)^3} \left(\frac{\partial f}{\partial t} - H p \frac{\partial f}{\partial p} \right) \quad \frac{dp}{dt} = -\left(\frac{\dot{a}}{a}\right) \cdot \frac{p_0}{a} = -H \cdot p(t)$$

$$\frac{d}{dt} \int \frac{d^3p}{(2\pi)^3} f = \frac{d}{dt} (n(t)) = \frac{dn}{dt}$$

$$-H \frac{1}{2\pi^2} \int_0^\infty dp p^3 \frac{\partial f}{\partial p} = 3H \frac{1}{2\pi^2} \left(\int dp p^2 f(p, t) \right) = 3Hn$$

integrate by parts

$$\frac{dn}{dt} + 3Hn = \frac{d}{dt} [f(p, t)]$$

if RHS = 0 $\rightarrow \frac{dn}{dt} = -3Hn \rightarrow \frac{dn}{n} = -3H dt = -3 \frac{1}{a} \frac{da}{dt} dt$

$$\frac{dn}{n} = -3 \frac{da}{a} \rightarrow \ln(n) = \ln(a^{-3}) + C = \log(c_1 \cdot a^{-3})$$

$$n(a) = n(a_1) \cdot \left(\frac{a_1}{a}\right)^3$$

if $RHS \neq 0$

$1+2 \leftrightarrow 3+4$

$$\frac{dn_i}{dt} + 3Hn_i = -A n_1 n_2 + B n_3 n_4$$

Hubble expansion destruction rate creation rate

in equilibrium: $A n_1^{eq} n_2^{eq} = B n_3^{eq} n_4^{eq}$

$$B = A \cdot \left(\frac{n_1^{eq} n_2^{eq}}{n_3^{eq} n_4^{eq}} \right)$$

$A \equiv \langle \sigma \cdot v \rangle$ thermally avg x-section x velocity

$$\frac{dn_i}{dt} + 3Hn_i = - \langle \sigma \cdot v \rangle \left(n_1 n_2 - \left(\frac{n_1^{eq} n_2^{eq}}{n_3^{eq} n_4^{eq}} \right) n_3 n_4 \right)$$

\hookrightarrow particle physics/QFT

$$n_i^{eq} = \int \frac{d^3p}{(2\pi)^3} \frac{1}{e^{\frac{-E_i - \mu_i}{T}} \pm 1}$$

Apply to $e^- + p^+ \leftrightarrow H + \gamma$

$$n_e = n_p$$

$$n_H \approx n_{H,eq}$$

$$n_{\gamma,eq} = \frac{2\zeta(3)}{\pi^2} T^3$$

$$X_e \equiv \frac{n_e}{n_p + n_H} = \frac{n_e}{n_b}$$

$$n_b = 2 n_{\gamma}^{eq}$$

$$\frac{1}{a^3} \cdot \frac{d(a^3 n_e)}{dt} = - \langle \sigma \cdot v \rangle \cdot \left(n_e n_p - \left(\frac{n_e^{eq} n_p^{eq}}{n_H^{eq} n_{\gamma}^{eq}} \right) n_H n_{\gamma} \right)$$

$n_e = n_p$

$$= - \langle \sigma \cdot v \rangle \left(n_e^2 - n_{e,eq}^2 \right)$$

free electron

$$n_e = n_p \approx n_b \cdot X_e$$

$$\text{b/c } n_b = n \cdot \frac{n_{\gamma}}{a^3}$$

$\rightarrow a^3 \cdot n_b = \text{const.}$ can bring outside

$$\frac{1}{a^3} \frac{d(n_e n_b a^3)}{dt} = \frac{n_b \cdot a^3}{a^3} \frac{dX_e}{dt} = n_b \cdot \frac{dX_e}{dt}$$

$$n_e = n_b \cdot X_e$$

$$n_b \cdot \frac{dX_e}{dt} = - \langle \sigma \cdot v \rangle \left(n_b^2 X_e^2 - (n_{b,eq})^2 (X_{e,eq})^2 \right)$$

Saha eqⁿ

$$\frac{dX_e}{dt} = - \langle \sigma \cdot v \rangle n_b \left(X_e^2 - X_{e,eq}^2 \right)$$

$\hookrightarrow n_b \cdot n_{\gamma} \propto T^3$

$T \propto \frac{1}{a(t)}$

$$\langle \sigma \cdot v \rangle \approx \sigma_T \cdot \left(\frac{B_H}{T} \right)^{3/2}$$

$$\frac{\sigma_T}{3} \frac{\alpha^2}{m_e^2}$$